

# Current fluctuations: mapping non-equilibrium to equilibrium

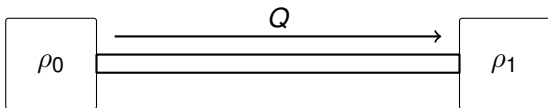
Alberto Imparato<sup>1</sup>, Vivien Lecomte<sup>2</sup>, Frédéric van Wijland<sup>3</sup>

<sup>1</sup>DPA, Aarhus <sup>2</sup>LPMA, Paris <sup>3</sup>MSC, Paris



GDR PHENIX – 24th October 2010

# Motivations



$$\text{Prob}(\mathcal{C}) \propto \exp \left\{ - \frac{\text{energy}(\mathcal{C})}{\text{temperature}} \right\} \quad \text{cannot describe}$$

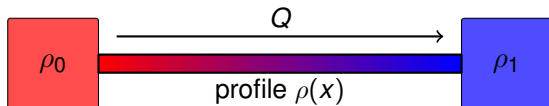
- Non-equilibrium steady-state

$$\text{Prob}[\rho(x)]$$

- Equilibrium fluctuations of dynamical observables

$$\text{Prob}[Q]$$

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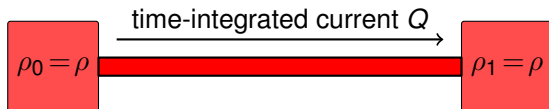
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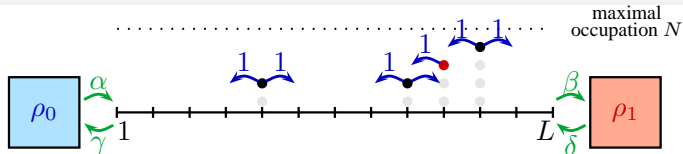
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# Exclusion processes



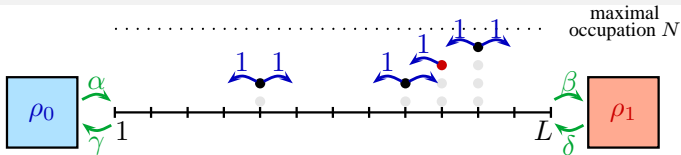
- Configurations: occupation numbers  $\{n_i\}$
- Exclusion rule:  $0 \leq n_i \leq N$
- Markov evolution

$$\partial_t P(\{n_i\}) = \sum_{n'_i} [W(n'_i \rightarrow n_i) P(\{n'_i\}) - W(n_i \rightarrow n'_i) P(\{n_i\})]$$

- Large deviation function of the time-integrated current  $Q$

$$\langle e^{-sQ} \rangle \sim e^{t\psi(s)} \quad (\Leftrightarrow \text{determining } P(Q))$$

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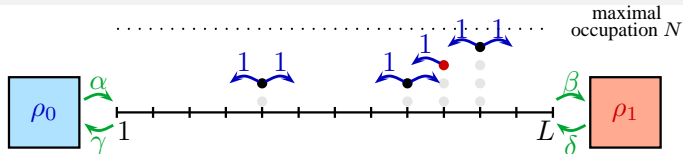
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## Example 1: distribution of current in the WASEP

- Periodic boundary condition (mean density  $\rho$ )
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- Periodic boundary condition (mean density  $\rho$ )
- Bulk field  $E$ : small asymmetry ( $\sim \frac{1}{L}$ ) in the jump rates  $\rightarrow$  non-eq.
- Distribution of the total current: (space&time integrated)

$$\frac{1}{t} \langle Q_{\text{tot}} \rangle = L E$$

$$\frac{1}{t} \langle Q_{\text{tot}}^2 \rangle_c = L 2\rho(1 - \rho)$$

$$\frac{1}{t} \langle Q_{\text{tot}}^4 \rangle_c = L^2 2[\rho(1 - \rho)]^2 \quad (\text{free particles would be } \sim L)$$

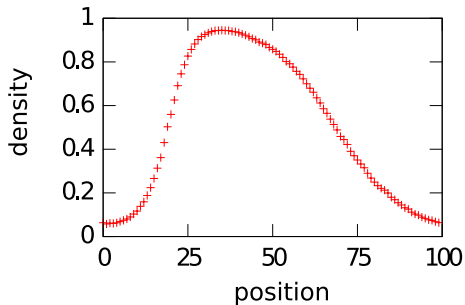
$$\vdots$$

$$\frac{1}{t} \langle Q_{\text{tot}}^k \rangle_c \sim L^{k-2} \quad (\text{anomalous scaling in } L)$$

$$\vdots$$

## Example 1: distribution of current in the WASEP

- Periodic boundary condition (mean density  $\rho$ )
- Bulk field  $E$ : small asymmetry ( $\sim \frac{1}{L}$ ) in the jump rates  $\rightarrow$  non-eq.
- Optimal profile leading to a large atypical flow  $J \equiv \frac{1}{t} \langle Q \rangle$ :
  - for  $|J - E/2| < J_c$ : flat profile
  - for  $|J - E/2| > J_c$ : heterogeneous profile

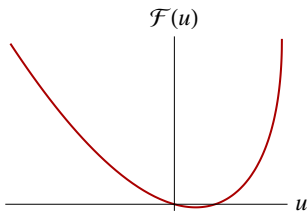


Non-steady and non-uniform density profile

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- Periodic boundary condition (mean density  $\rho$ )
- Bulk field  $E$ : small asymmetry ( $\sim \frac{1}{L}$ ) in the jump rates  $\rightarrow$  non-eq.
- With a field [Appert, Derrida, VL, van Wijland, PRE **78** 021122 (2008)]

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{2} \mathbf{s}(\mathbf{s} - E) \frac{\langle Q^2 \rangle_c}{t}}_{\text{at saddle-point}} + \underbrace{L^{-2} D\mathcal{F}(u)}_{\substack{\text{small fluctuations} \\ \text{(determinant)}}} \quad \text{with } u = \underbrace{-\mathbf{s}(\mathbf{s} - E) \frac{\sigma\sigma''}{16D^2}}_{\text{can become } > 0}$$

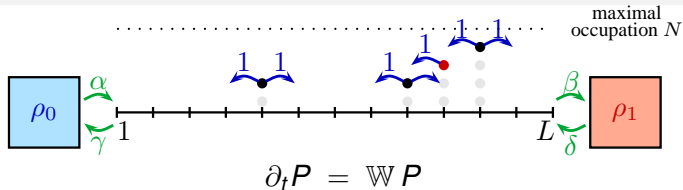


Dynamical phase transition  
between  
stationary and non-stationary  
profiles

What about  
**boundary** driven  
non-equilibrium transport?

## Operator representation

[Schütz &amp; Sandow PRE 49 2726]



$$\mathbb{W} = \sum_{1 \leq k \leq L-1} [S_k^+ S_{k+1}^- + S_k^- S_{k+1}^+ - \hat{n}_k (1 - \hat{n}_{k+1}) - \hat{n}_{k+1} (1 - \hat{n}_k)]$$

$$+ \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1]$$

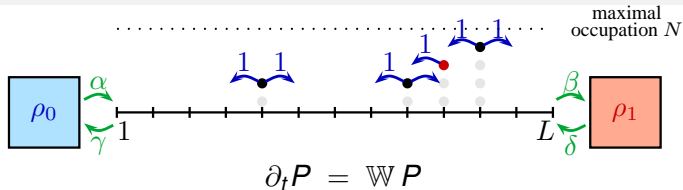
$$+ \delta [S_L^+ - (1 - \hat{n}_L)] + \beta [S_L^- - \hat{n}_L]$$

$S^\pm$  and **creation and annihilation operators:**

$$S^+ |n\rangle = (N - n) |n + 1\rangle \quad S^- |n\rangle = n |n - 1\rangle$$

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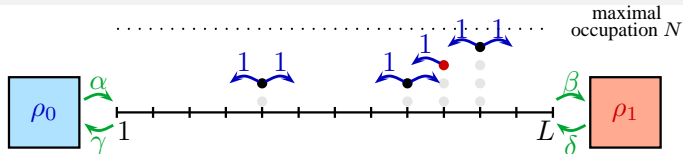
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$S^\pm$  and  $S^z = \hat{n} - \frac{N}{2}$  are **spin operators** (with  $j = \frac{N}{2}$ )

# Large deviation function



$$\langle e^{-sQ} \rangle \sim e^{t\psi(s)} \quad \text{with} \quad \psi(s) = \max \text{Sp } \mathbb{W}(s)$$

$$\begin{aligned} \mathbb{W}(s) = & \sum_{1 \leq k \leq L-1} [S_k^+ S_{k+1}^- + S_k^- S_{k+1}^+ - \hat{n}_k(1 - \hat{n}_{k+1}) - \hat{n}_{k+1}(1 - \hat{n}_k)] \\ & + \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1] \\ & + \delta [S_L^+ e^s - (1 - \hat{n}_L)] + \beta [S_L^- e^{-s} - \hat{n}_L] \end{aligned}$$

# Mapping non-eq to eq [Imparato, VL, van Wijland, PTP **184** 276 (2010)]

## Large deviations of the current

$$\psi(\mathbf{s}) = \max_{\text{Sp}} \mathbb{W}(\mathbf{s})$$

$$\mathbb{W}(\mathbf{s}) = \overbrace{\sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1}}^{\text{invariant by rotation}} + \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1] + \delta [S_L^+ e^{\mathbf{s}} - (1 - \hat{n}_L)] + \beta [S_L^- e^{-\mathbf{s}} - \hat{n}_L]$$



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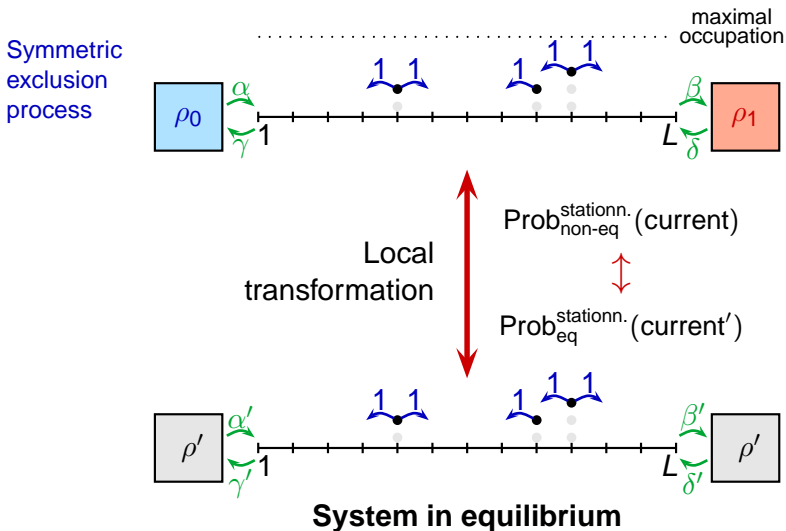
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## Local transformation

$$\mathcal{Q}^{-1} \mathbb{W}(\mathbf{s}) \mathcal{Q} = \sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1} + \alpha' [S_1^+ - (1 - \hat{n}_1)] + \gamma' [S_1^- - \hat{n}_1] + \delta' [S_L^+ e^{\mathbf{s}'} - (1 - \hat{n}_L)] + \beta' [S_L^- e^{-\mathbf{s}'} - \hat{n}_L]$$

describes contact with reservoirs of same densities

# Summary [Imparato, VL, van Wijland, PRE **80** 011131 & PTP **184** 276 (2010)]



# Result

[Imparato, VL, van Wijland, PRE **80** 011131 & PTP **184** 276 (2010)]

- Large deviation function

[Reminder:  $\langle e^{-sQ} \rangle \sim e^{t\psi(s)}$ ]

$$\psi(s) = \underbrace{\frac{1}{L}\mu(s)}_{\text{saddle}} + \underbrace{\frac{D}{8L^2}\mathcal{F}\left(\frac{\sigma''}{2D^2}\mu(s)\right)}_{\text{fluctuations}}$$

$$\mu(s) = -(\operatorname{arcsinh} \sqrt{\omega})^2 \quad \omega = (1 - e^s)(e^{-s}\rho_0 - \rho_1 - (e^{-s} - 1)\rho_0\rho_1)$$

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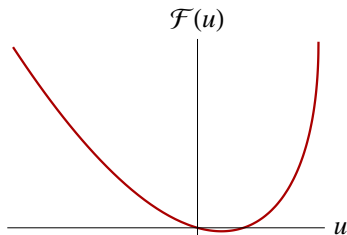
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- Universal function  $\mathcal{F}$

$$\mathcal{F}(u) = \sum_{k \geq 2} \frac{(-2u)^k B_{2k-2}}{\Gamma(k)\Gamma(k+1)}$$



# Summary

## Approach:

- large deviation function
- operator formalism
- fluctuating hydrodynamics

## Results:

- universal fluctuations
- possible dynamical phase transition
- non-eq. to eq. mapping

## Open questions:

- Eq $\leftrightarrow$ non-eq mapping in higher dimensions?
- More generic systems of interacting particles?
- Link to other eq $\leftrightarrow$ non-eq mappings?
- Crossover to KPZ? Universal fluctuations?

# Finite-size effects

[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

## For **periodic boundary conditions**

- large deviation function

$$\psi(\mathbf{s}) = \underbrace{L^{-1} \rho(1 - \rho) \mathbf{s}^2}_{\text{order 0}} + \underbrace{L^{-2} \mathcal{F}(u)}_{\text{finite-size}} \quad \text{with} \quad u = -\frac{1}{2} \rho(1 - \rho) \mathbf{s}^2$$

- universal function

$$\mathcal{F}(u) = \sum_{k \geq 2} \frac{(-2u)^k B_{2k-2}}{\Gamma(k) \Gamma(k+1)}$$

- scaling of cumulants of the total current  $Q_{\text{tot}}$

$$\frac{1}{t} \langle Q_{\text{tot}}^2 \rangle \sim L$$

$$\frac{1}{t} \langle Q_{\text{tot}}^{2k} \rangle \sim L^{2k-2} \quad (k \geq 2)$$

# Finite-size effects

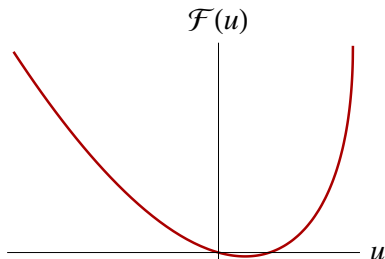
[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

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# With a field

[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

Periodic boundary conditions

Driving field  $E$



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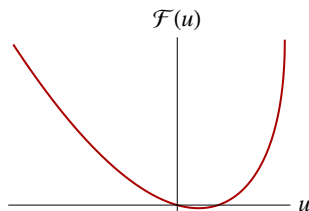
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Periodic boundary conditions  
Driving field  $E$

For the WASEP:  
 $D = 1, \sigma = 2\rho(1 - \rho)$

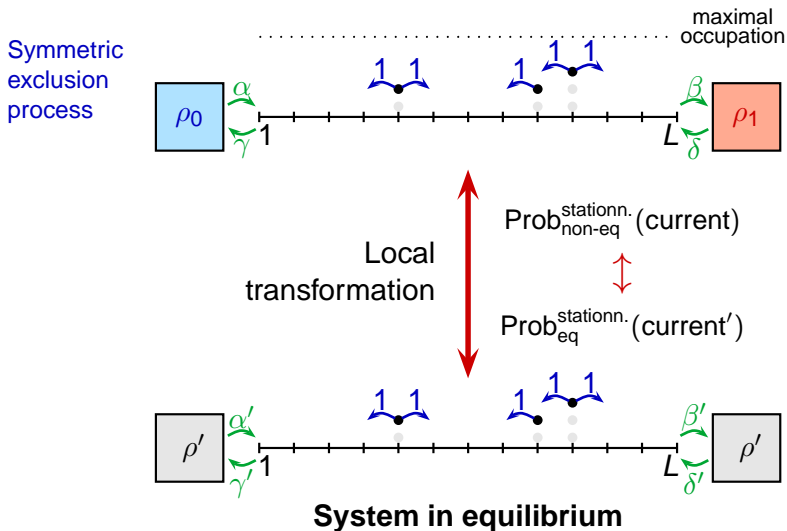
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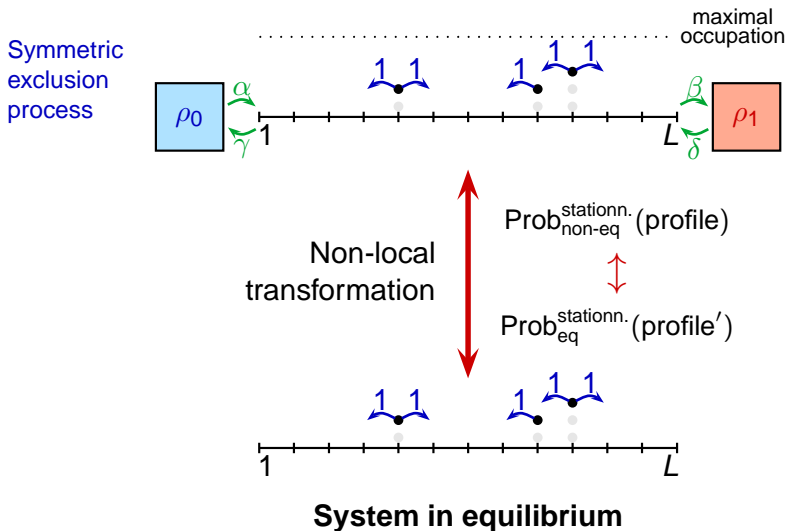
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## For the current

[Imparato, VL, van Wijland, **PRE** 80 011131]

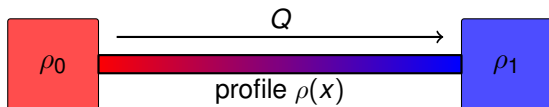
## For the density profile

[Tailleur, Kurchan, VL, JPA 41 505001]



# Non-local mapping

[Tailleur, Kurchan, VL, JPA **41** 505001]

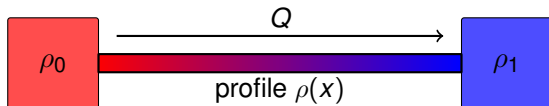


## Boundary-driven transport model:

- long-range correlations
- breaking of time-reversal symmetry

# Non-local mapping

[Tailleur, Kurchan, VL, JPA **41** 505001]



## Boundary-driven transport model:

- long-range correlations
- breaking of time-reversal symmetry

## Non-local mapping to equilibrium:

- accounts for long-range correlations
- (density gradient)<sub>non-eq.</sub>  $\longleftrightarrow$  (fixed density)<sub>eq.</sub>
- yields  $\text{Prob}[\rho(x)]$  through an extremalization principle