Motion of interfaces with an internal degree of freedom: simple models and open questions

Stewart Barnes[1,2], Jean-Pierre Eckmann[3], Thierry Giamarchi[1]

[1]DQMP, Genève, Suisse
[2]Physics Department, University of Miami, USA
[3]Département de Physique Théorique et Section de Mathématiques, Genève

Genève – 14 November 2019

Motion of interfaces with an internal degree of freedom: simple models and open questions

Stewart Barnes[1,2], Jean-Pierre Eckmann[3], Thierry Giamarchi[1]
Alejandro Kolton[4], Victor Purrello[4,5]

[1]DQMP, Genève, Suisse
[2]Physics Department, University of Miami, USA
[3]Département de Physique Théorique et Section de Mathématiques, Genève
[4]CNEA, Bariloche, Argentina [5]IFIMAR, Mar del Plata, Argentina

Genève – 14 November 2019

Outline

1. Examples
   - Systems
   - Questions
2. A panorama of results
   - Zero temperature
   - Small temperature
3. Role of hidden degrees of freedom
   - Examples
   - Results
4. Open questions and perspectives
   - Effective inertia
   - Active/soft matter

Outline

Outline

1. Examples
   - Systems
   - Questions
2. A panorama of results
   - Zero temperature
   - Small temperature
3. Role of hidden degrees of freedom
   - Examples
   - Results
4. Open questions and perspectives
   - Effective inertia
   - Active/soft matter

Outline

1. Examples
   - Systems
   - Questions
2. A panorama of results
   - Zero temperature
   - Small temperature
3. Role of hidden degrees of freedom
   - Examples
   - Results
4. Open questions and perspectives
   - Effective inertia
   - Active/soft matter

T. Halpin-Healy, Y.-C. Zhang/Physics Reports 254 (1995) 215–414

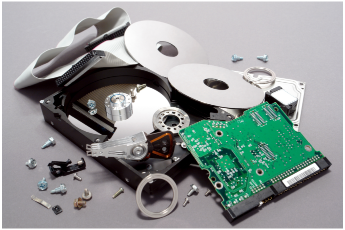
1. Introduction

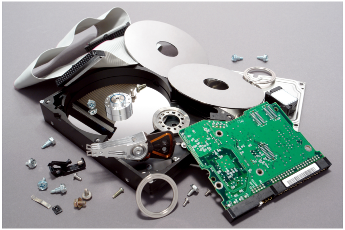
Consider a wheat field of dark golden hue and densely planted in level ground, being roughly rectangular in shape, but rather large in extent, and stretching lazily toward the distant horizon. On a cool, but calm August evening, with nary a breeze about, the edge of the field is ignited, in preparation for leaving the soil fallow the following season. The propagating fire front, initially straight by virtue of its birth along the edge, evolves in a kinetic, violent fashion and heads mercilously into the bulk of the field. Burning shafts of wheat communicate the conflagaration locally to their neighbors, and the narrow, bright, and tortuously shaped fire line, an interface separating the blackened region from the portion of the field soon to be consumed, becomes increasingly rough as random elements, such as local inhomogeneities in the moisture content or density of the wheat, begin to have a large scale cumulative effect.

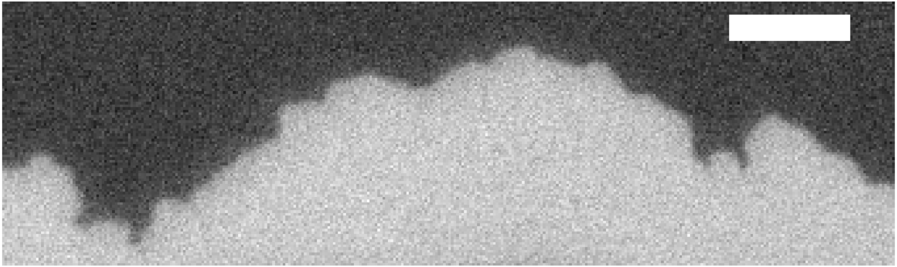
T. Halpin-Healy, Y.-C. Zhang / Physics Reports 254 (1995) 215–414

1. Introduction

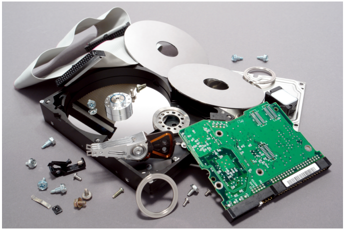
Consider a wheat field of dark golden hue and densely planted in level ground, being roughly rectangular in shape, but rather large in extent, and stretching lazily toward the distant horizon. On a cool, but calm August evening, with nary a breeze about, the edge of the field is ignited, in preparation for leaving the soil fallow the following season. The propagating fire front, initially straight by virtue of its birth along the edge, evolves in a kinetic, violent fashion and heads mercilously into the bulk of the field. Burning shafts of wheat communicate the conflagaration locally to their neighbors, and the narrow, bright, and tortuously shaped fire line, an interface separating the blackened region from the portion of the field soon to be consumed, becomes increasingly rough as random elements, such as local inhomogeneities in the moisture content or density of the wheat, begin to have a large scale cumulative effect.

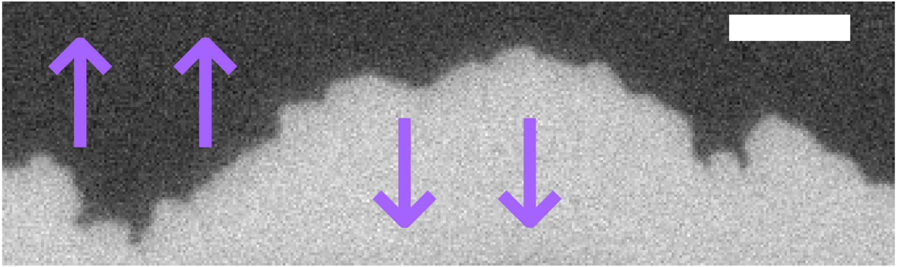




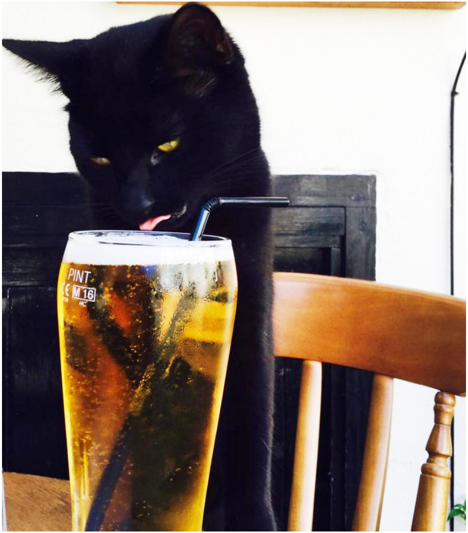


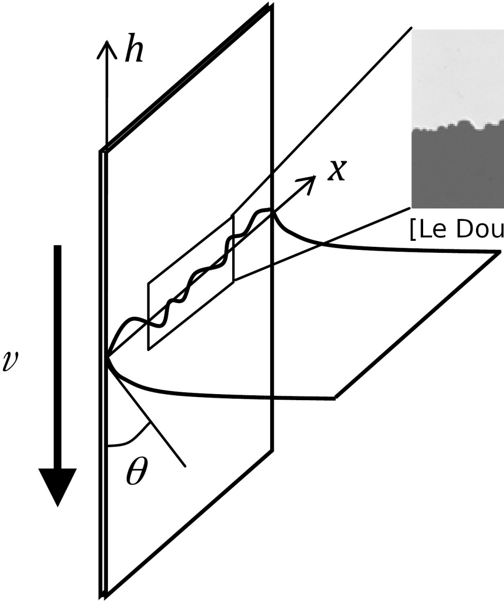
[Metaxas *et al.* PRL **99** 217208 (2007)]

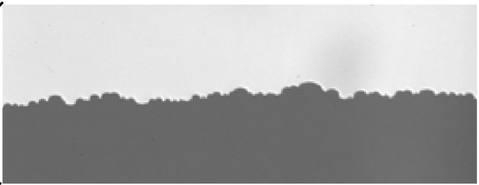




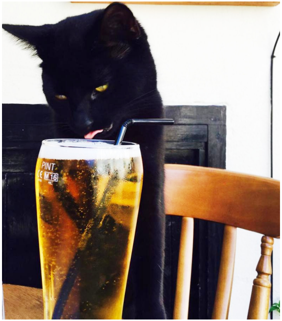
[Metaxas *et al.* PRL **99** 217208 (2007)]


PINT
M 16



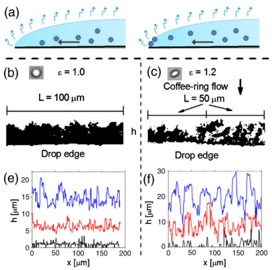


[Le Doussal, Wiese, Moulinet, Rolley,
EPL **87** 56001 (2009)]



Front of evaporation/imbibition

PRL **110,** 035501 (2013)



PRL **110,** 035501 (2013) PHYSICAL REVIEW LETTERS week ending 18 JANUARY 2013

Effects of Particle Shape on Growth Dynamics at Edges of Evaporating Drops of Colloidal Suspensions

Peter J. Yunker,[1] Matthew A. Lohr,[1] Tim Still,[1,2] Alexei Borodin,[3] D. J. Durian,[1] and A. G. Yodh[1]

[1]*Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*
[2]*Complex Assemblies of Soft Matter, CNRS-Rhodia-University of Pennsylvania, UMI 3254, Bristol, Pennsylvania 19007, USA*
[3]*Department of Mathematics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*

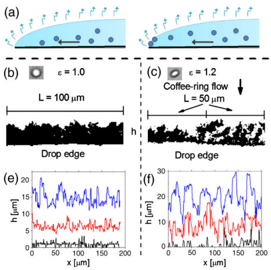
(Received 23 July 2012; published 18 January 2013)

We study the influence of particle shape on growth processes at the edges of evaporating drops. Aqueous suspensions of colloidal particles evaporate on glass slides, and convective flows during evaporation carry particles from drop center to drop edge, where they accumulate. The resulting particle deposits grow inhomogeneously from the edge in two dimensions, and the deposition front, or growth line, varies spatiotemporally. Measurements of the fluctuations of the deposition front during evaporation enable us to identify distinct growth processes that depend strongly on particle shape. Sphere deposition exhibits a classic Poisson-like growth process; deposition of slightly anisotropic particles, however, belongs to the Kardar-Parisi-Zhang universality class, and deposition of highly anisotropic ellipsoids appears to belong to a third universality class, characterized by Kardar-Parisi-Zhang fluctuations in the presence of quenched disorder.

DOI: 10.1103/PhysRevLett.110.035501 PACS numbers: 61.43.Fs, 64.70.kj, 64.70.pv, 82.70.Dd

Front of evaporation/imbibition

PRL **110,** 035501 (2013)



PRL **110,** 035501 (2013) PHYSICAL REVIEW LETTERS week ending 18 JANUARY 2013

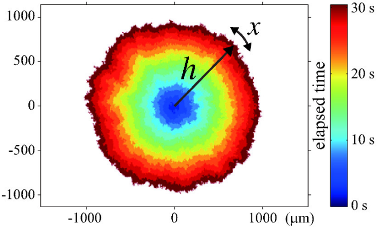
Effects of Particle Shape on Growth Dynamics at Edges of Evaporating Drops of Colloidal Suspensions

Peter J. Yunker,[1] Matthew A. Lohr,[1] Tim Still,[1,2] Alexei Borodin,[3] D. J. Durian,[1] and A. G. Yodh[1]
[1]*Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*
[2]*Complex Assemblies of Soft Matter, CNRS-Rhodia-University of Pennsylvania, UMI 3254, Bristol, Pennsylvania 19007, USA*
[3]*Department of Mathematics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA*
(Received 23 July 2012; published 18 January 2013)

We study the influence of particle shape on growth processes at the edges of evaporating drops. Aqueous suspensions of colloidal particles evaporate on glass slides, and convective flows during evaporation carry particles from drop center to drop edge, where they accumulate. The resulting particle deposits grow inhomogeneously from the edge in two dimensions, and the deposition front, or growth line, varies spatiotemporally. Measurements of the fluctuations of the deposition front during evaporation enable us to identify distinct growth processes that depend strongly on particle shape. Sphere deposition exhibits a classic Poisson-like growth process; deposition of slightly anisotropic particles, however, belongs to the Kardar-Parisi-Zhang universality class, and deposition of highly anisotropic ellipsoids appears to belong to a third universality class, characterized by Kardar-Parisi-Zhang fluctuations in the presence of quenched disorder.

DOI: 10.1103/PhysRevLett.110.035501 PACS numbers: 61.43.Fs, 64.70.kj, 64.70.pv, 82.70.Dd

Front between turbulent modes in liquid crystals



J Stat Phys (2015) 160:794–814
DOI 10.1007/s10955-015-1282-1

A KPZ Cocktail-Shaken, not Stirred...

Toasting 30 Years of Kinetically Roughened Surfaces

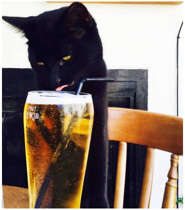
Timothy Halpin-Healy[1] · Kazumasa A. Takeuchi[2,3]

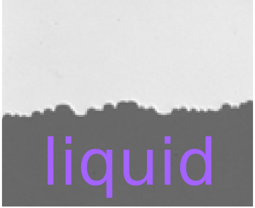
Large range of physical scales

Wide spectrum of phenomena

Questions:

- What are the relevant physical features?
- How to characterise the geometry and the dynamics?



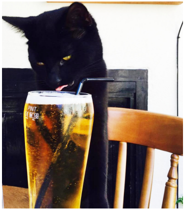


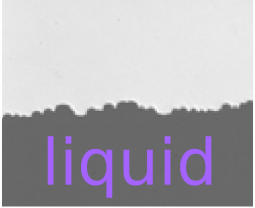
Order: surface tension

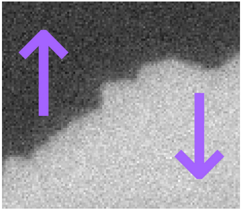
Disorder: substrate impurities

Drive: imposed $\langle$liquid level$\rangle$

Noise: negligible







Order: surface tension

Disorder: substrate impurities

Drive: imposed ⟨liquid level⟩

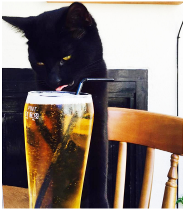
Noise: negligible

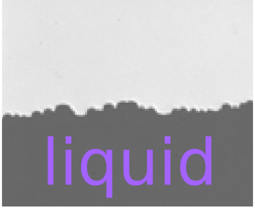
Order: energetic cost of interface

Disorder: (magnetic) impurities

Drive: external magnetic field

Noise: thermal



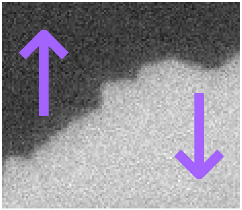


Order: surface tension

Disorder: substrate impurities

Drive: imposed ⟨liquid level⟩

Noise: negligible



Order: energetic cost of interface

Disorder: (magnetic) impurities

Drive: external magnetic field

Noise: thermal

Order: convex patches burn faster

Disorder: inhomogeneities

Drive: instability of unburnt grass

Noise: turbulence in the air

Large range of physical scales

Wide spectrum of phenomena

Questions:

- What are the relevant physical features?

- How to characterise the geometry and the dynamics?

Large range of physical scales

Wide spectrum of phenomena

Questions:

- What are the relevant physical features?
  - Competition between
    order (tends to align)
    quenched **disorder** (tends to deform)
  - **Noise** (space and time fluctuating force)
  - **Drive** (external force or internal instability)

- How to characterise the geometry and the dynamics?

Large range of physical scales

Wide spectrum of phenomena

Questions:

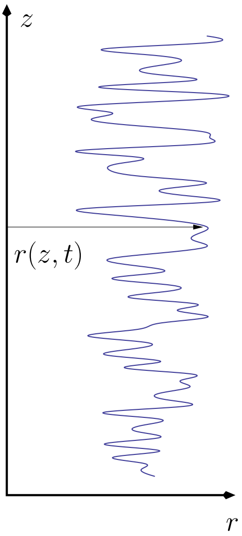
- What are the relevant physical features?
  - Competition between
    order (tends to align)
    quenched **disorder** (tends to deform)
  - **Noise** (space and time fluctuating force)
  - **Drive** (external force or internal instability)
- When is disorder *relevant*?
- How to characterise the geometry and the dynamics?

Outline

1. Examples
   - Systems
   - Questions
2. A panorama of results
   - Zero temperature
   - Small temperature

Outline

1. Examples
   - Systems
   - Questions
2. A panorama of results
   - Zero temperature
   - Small temperature



z
$r(z,t)$
r

- **Geometry**:
  - Roughness function B

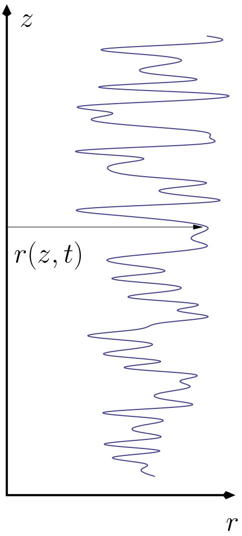
$$B(z,t) = \overline{\left\langle \left[r(z_1+z,t_1+t) - r(z_1,t_1) \right]^2 \right\rangle}$$

$\langle \ldots \rangle$ = thermal average

$\overline{\ldots}$ = disorder average

  - Roughness exponent ζ

$$B(z,0) \sim z^{2\zeta} \qquad (z \to \infty)$$



- **Geometry**:
  - Roughness function B

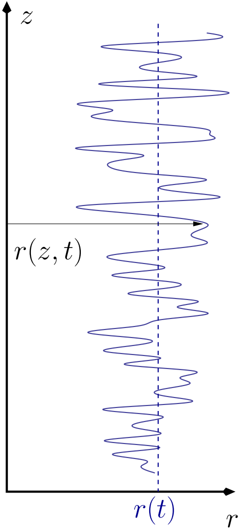
$$B(z,t) = \overline{\left\langle \left[r(z_1+z, t_1+t) - r(z_1,t_1) \right]^2 \right\rangle}$$

$\langle \dots \rangle$ = thermal average

$\overline{\dots}$ = disorder average

  - Roughness exponent ζ

$$B(z,0) \sim z^{2\zeta} \qquad (z \to \infty)$$



- **Geometry**:
  - Roughness function B

$$B(z,t) = \overline{\langle [r(z_1+z,t_1+t) - r(z_1,t_1)]^2 \rangle}$$

$\langle \ldots \rangle$ = thermal average

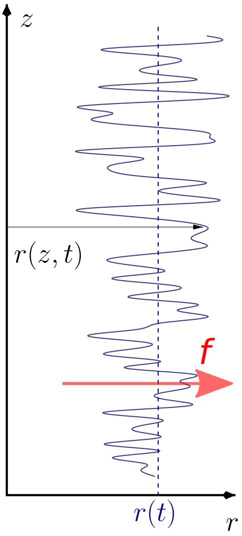
$\overline{\ldots}$ = disorder average

  - Roughness exponent ζ

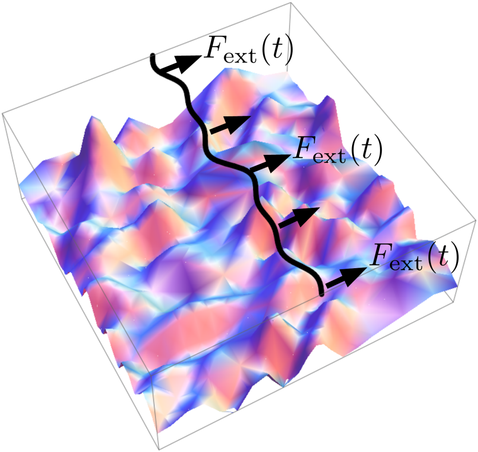
$$B(z,0) \sim z^{2\zeta} \qquad (z \to \infty)$$

- **Dynamics**:
  - Velocity-force characteristic

$$v(f) = \overline{\langle \partial_t r(t) \rangle} \qquad (t \to \infty)$$



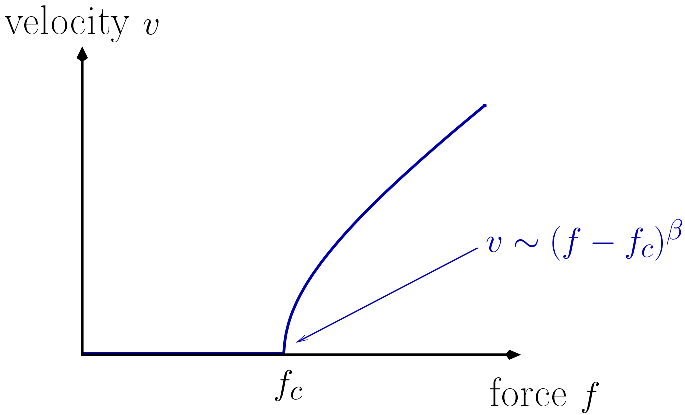
The velocity-force characteristic $v(f)$



[E. Agoritsas]

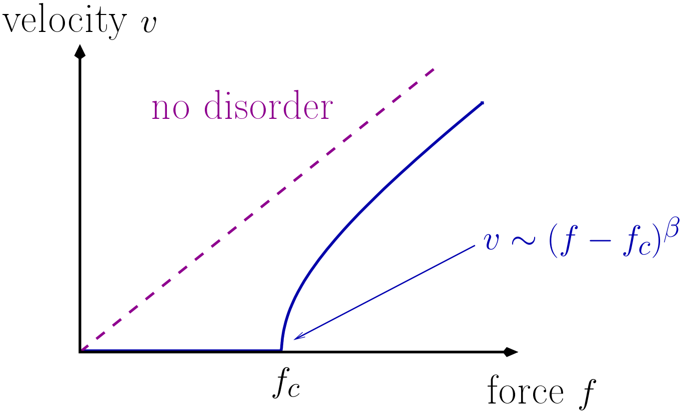
Schematic representation of the protocol.

Depinning transition @ zero temperature



Criticality at a threshold force f_c [*non-equilibrium* phase transition]

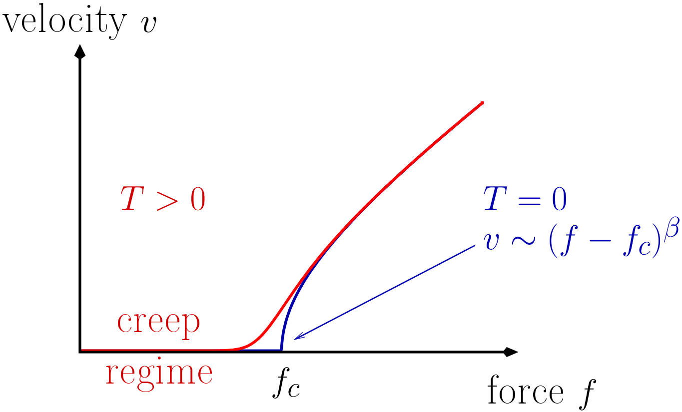
Depinning transition @ zero temperature



Criticality at a threshold force f_c [*non-equilibrium* phase transition]

Disorder is **relevant**

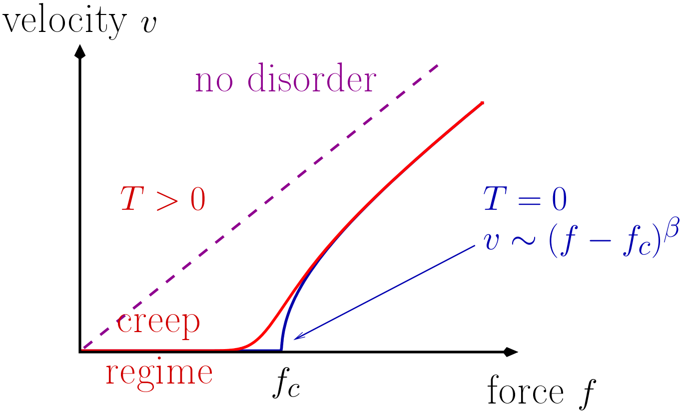
Depinning transition & finite temperature



Creep law: $v(f) \overset{f\to 0}{\sim} e^{-\frac{U_c}{T}(f_c/f)^{\mu}}$

[**Highly non-linear** response]

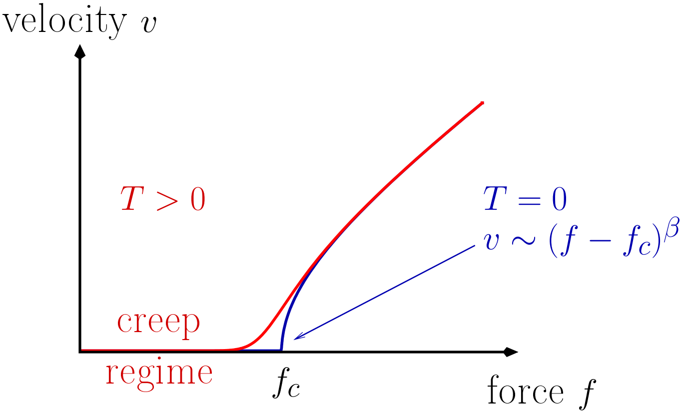
Depinning transition & finite temperature



Creep law: $v(f) \overset{f \to 0}{\sim} e^{-\frac{U_c}{T}(f_c/f)^\mu}$ **[Highly non-linear** response**]**

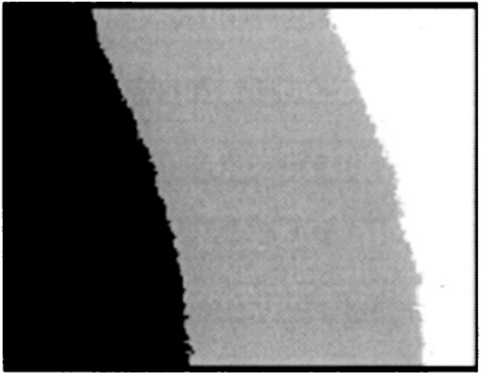
Disorder is **relevant**

Depinning transition & finite temperature



Creep law: $v(f) \overset{f \to 0}{\sim} e^{-\frac{U_c}{T}(f_c/f)^{\mu}}$ [**Highly non-linear** response]

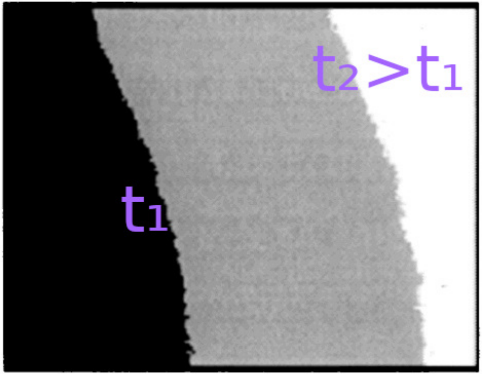
$$v(f) \sim \exp\left[-\frac{U_c}{T}\left(\frac{f_c}{f}\right)^{\mu}\right] \qquad \text{(creep law)}$$



$\mu = 1/4$ here

[Lemerle *et al.*, PRL **80** 849 (1998)]

$$v(f) \sim \exp\left[-\frac{U_c}{T}\left(\frac{f_c}{f}\right)^{\mu}\right]$$ (creep law)

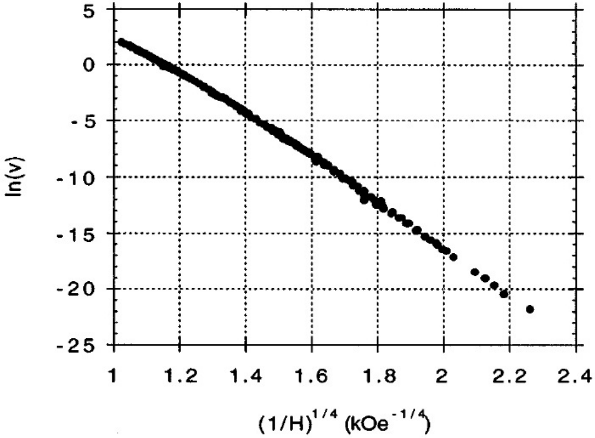


$\mu = 1/4$ here

[Lemerle *et al.*, PRL **80** 849 (1998)]

$$v(f) \sim \exp\left[-\frac{U_c}{T}\left(\frac{f_c}{f}\right)^{\mu}\right]$$

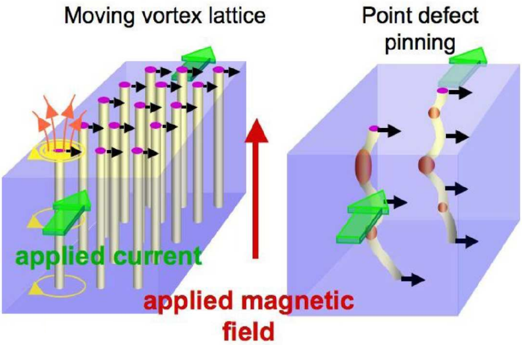
(creep law)



$\mu = 1/4$ here

[Lemerle *et al.*, PRL **80** 849 (1998)]

A remarkable consequence



[Matti Irjala]

Without disorder, type II superconductors would dissipate Ohmically.

A simple picture of depinning?

Disordered elastic systems

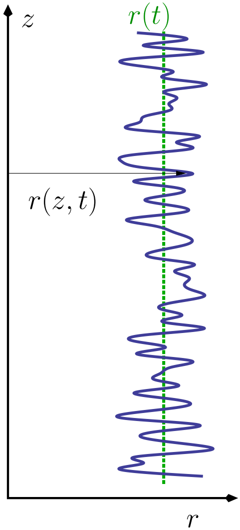
- Elasticity: tends to flatten the interface

$$\mathcal{H}^{\text{el}} = \frac{c}{2} \int dz \, \left(\nabla r(z)\right)^2 \qquad \textbf{[Short-range]}$$

$$\mathcal{H}^{\text{el}} = \frac{c}{2\pi} \int dz dz' \, \frac{\left(r(z) - r(z')\right)^2}{(z - z')^2} \qquad \text{[Long-range]}$$

- Disorder: tends to deform it

$$\mathcal{H}_V^{\text{dis}} = \int dz \; V\left(z, r(z)\right)$$



Competition btw “order” and “disorder”

Disordered elastic systems

- Elasticity: tends to flatten the interface

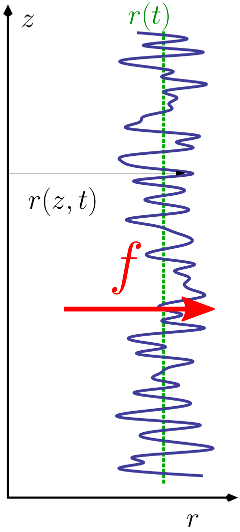
$$\mathcal{H}^{\text{el}} = \frac{c}{2}\int dz\, \left(\nabla r(z)\right)^2 \qquad \textbf{[Short-range]}$$

$$\mathcal{H}^{\text{el}} = \frac{c}{2\pi}\int dz dz' \,\frac{\left(r(z)-r(z')\right)^2}{(z-z')^2} \qquad \text{[Long-range]}$$

- Disorder: tends to deform it

$$\mathcal{H}_V^{\text{dis}} = \int dz\; V\big(z, r(z)\big)$$

- Force: can induce a motion of the interface



Competition btw "order" and "disorder"

Some known results

1. Huge variety of physical systems and theoretical approaches
   - ⋆ Elastic manifolds (lines, membranes, interfaces); periodic (vortex lattices); growth interfaces (aggregation, wetting)
   - ⋆ Methods: field theory, renormalisation group, scaling analysis, exactly solvable models, “replica trick”
   - ⋆ Reviews: Halpin-Healy&Zhang; Blatter&*al.*; Quastel; Corwin
2. Nature of fluctuations in **dimension 1+1** (elastic line)
   - ⋆ No disorder ($V(z,r) \equiv 0$):
     diffusive ($r \sim z^{1/2}$), **Edwards-Wilkinson** (EW)
   - ⋆ Disorder ($V(z,r) \not\equiv 0$):
     super-diffusive ($r \sim z^{2/3}$), **Kardar-Parisi-Zhang** (KPZ)

Some known results

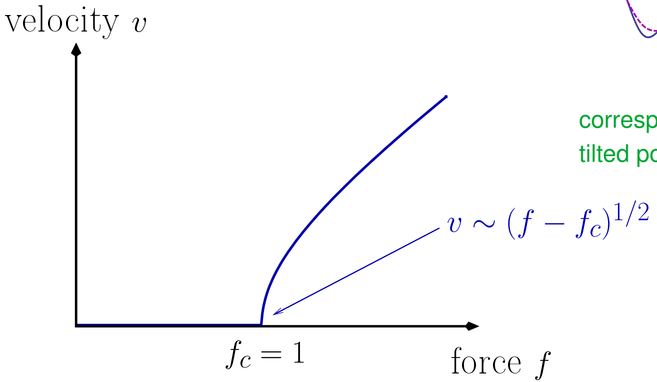
1. Huge variety of physical systems and theoretical approaches
   - ⋆ Elastic manifolds (lines, membranes, interfaces); periodic (vortex lattices); growth interfaces (aggregation, wetting)
   - ⋆ Methods: field theory, renormalisation group, scaling analysis, exactly solvable models, “replica trick”
   - ⋆ Reviews: Halpin-Healy&Zhang; Blatter&*al.*; Quastel; Corwin
2. Nature of fluctuations in **dimension 1+1** (elastic line)
   - ⋆ No disorder ($V(z, r) \equiv 0$):
     diffusive ($r \sim z^{1/2}$), **Edwards-Wilkinson** (EW)
   - ⋆ Disorder ($V(z, r) \not\equiv 0$):
     super-diffusive ($r \sim z^{2/3}$), **Kardar-Parisi-Zhang** (KPZ)

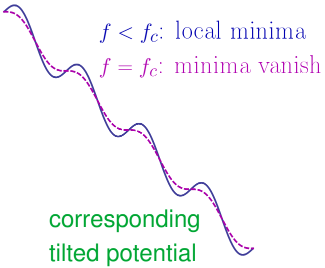
   Disorder is **always** relevant

Depinning @ zero temperature

Effective model for the mean interface position $r(t)$

$$\alpha \partial_t r = \underbrace{f - \cos \kappa r}_{\text{total force}}$$

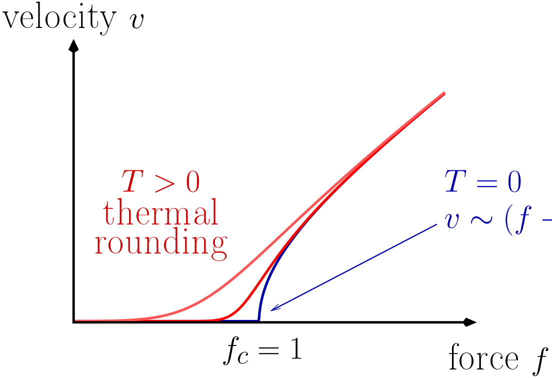


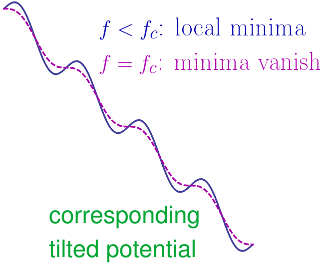


Depinning @ finite temperature

Effective model for the mean interface position $r(t)$

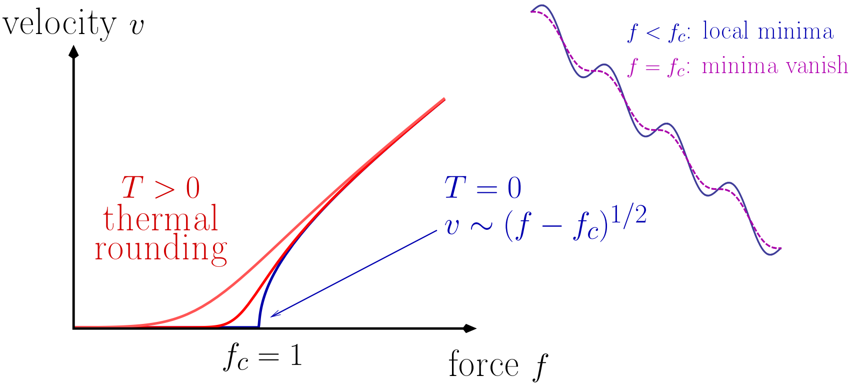
$$\alpha\partial_t r = \underbrace{f - \cos\kappa r}_{\text{total force}} + \eta$$





Depinning @ finite temperature

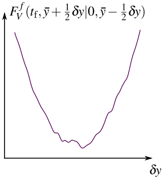
Effective model for the mean interface position *r*(*t*)

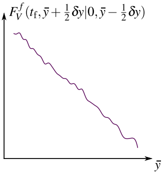


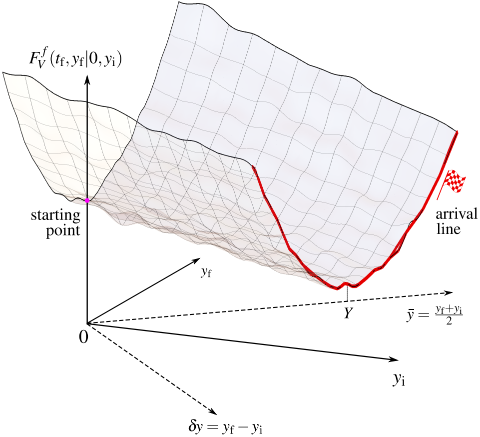
- Depinning: **ok**, but “mean-field” depinning exponent $1/2$
- Creep: **not ok** (linear response at $f \ll f_c$ instead of creep law)

A simple picture of creep?

From one to **two** effective degrees of freedom







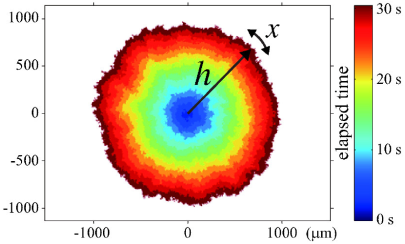
[Elisabeth Agoritsas, Reinaldo García-García, VL, Lev Truskinovsky and Damien Vandembroucq, J. Stat. Phys. **164** 1394 (2016)]

Geometry: the roughness $B(z)$

Roughness function $B(z)$ (now at equal times)

$$B(z) = \overline{\left\langle \left[r(z_1 + z, t) - r(z_1, t) \right]^2 \right\rangle} \sim z^{2\zeta} \qquad (z \to \infty)$$

Front between turbulent modes in liquid crystals



J Stat Phys (2015) 160:794–814
DOI 10.1007/s10955-015-1282-1

A KPZ Cocktail-Shaken, not Stirred...

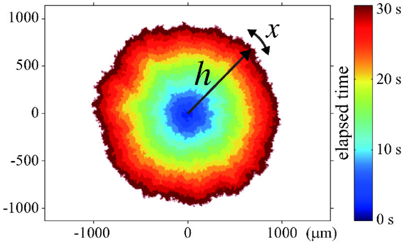
Toasting 30 Years of Kinetically Roughened Surfaces

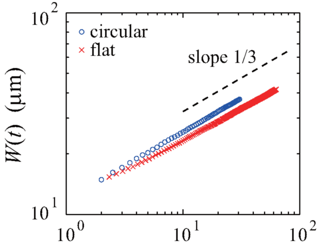
Timothy Halpin-Healy[1] · Kazumasa A. Takeuchi[2,3]

Roughness function $B(z)$ (now at equal times)

$$B(z) = \overline{\left\langle \left[r(z_1 + z, t) - r(z_1, t) \right]^2 \right\rangle} \sim z^{2\zeta} \qquad (z \to \infty)$$

Front between turbulent modes in liquid crystals

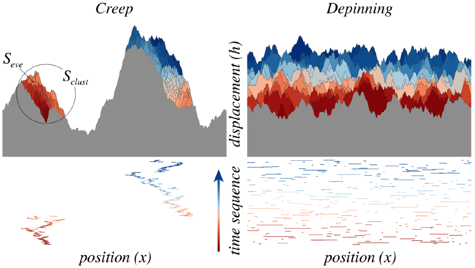


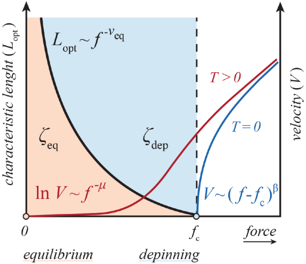


Roughness function $B(z)$ (now at equal times)

$$B(z) = \overline{\left\langle \left[r(z_1 + z, t) - r(z_1, t) \right]^2 \right\rangle} \sim z^{2\zeta} \qquad (z \to \infty)$$

At non-zero f:

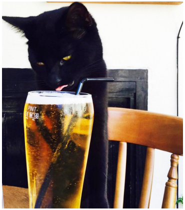


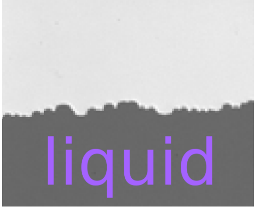


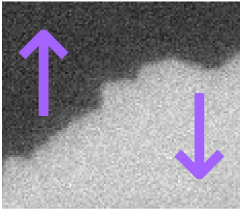
[Ezequiel Ferrero, Laura Foini, Thierry Giamarchi, Alejandro Kolton, Alberto Rosso,PRL **118** 147208 (2017)]

As f increases, ζ moves from $\zeta_{eq} = \frac{2}{3}$ to $\zeta_{dep} \approx 1.15$

Outline



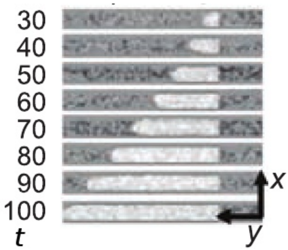

liquid



Is the knowledge of $r(z)$ sufficient?

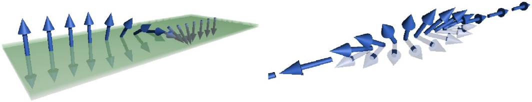
$\rightarrow$ Consider the dynamics of simple examples.

Spintronics

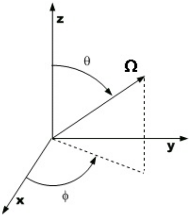


from Yamanouchi *et al.*, Science **317** 1726 (2007)

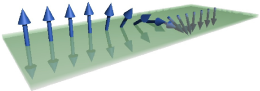
Bulk model $\rightsquigarrow$ effective description

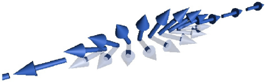


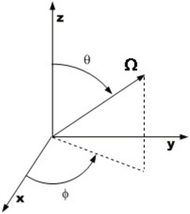
from Tatara *et al.*, J. Phys. Soc. Jap **77** 031003 (2008)



Bulk model $\rightsquigarrow$ effective description







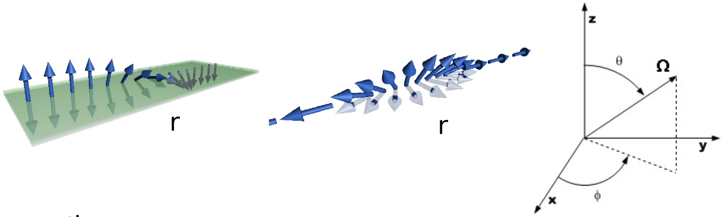
- Bulk energy

$$E = \int d^d x \left\{ J \left[(\nabla\theta)^2 + \sin^2\theta (\nabla\phi)^2 \right] + K \sin^2\theta + K_\perp \sin^2\theta \cos^2\phi \right\}$$

- Equation of motion (Landau-Lifshitz-Gilbert)

$$\partial_t \Omega = \Omega \times \left(\frac{\delta E}{\delta \Omega} + f + \eta \right) - \Omega \times \left(\alpha \partial_t \Omega \right)$$

Bulk model $\rightsquigarrow$ effective description

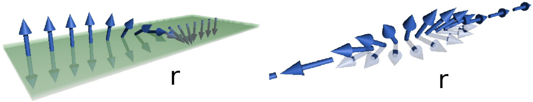


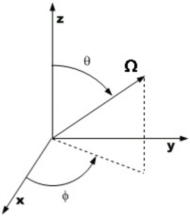
- Effective equations

$$\begin{aligned}
\alpha\partial_t r - \partial_t\phi &= J(\nabla r)^2 + F_{\text{pinning}} + f_{\text{ext}} + \eta_1 \\
\underbrace{\alpha\partial_t\phi + \partial_t r}_{\text{damping}} &= \underbrace{J(\nabla\phi)^2 - \frac{1}{2}K_\perp \sin 2\phi}_{\text{forces}} \; \underbrace{+\,\eta_2}_{\text{thermal noise}}
\end{aligned}$$

- Further simplification:
  Model reduction (from many to few degrees of freedom)

Bulk model $\rightsquigarrow$ effective description





- Case of a rigid wall (no elasticity)

$$\alpha\partial_t r - \partial_t\phi = \overbrace{-\cos\kappa r}^{\text{pinning}} + \overbrace{f}^{\text{external}} + \eta_1$$

$$\alpha\partial_t\phi + \partial_t r = -\frac{1}{2}K_\perp \sin 2\phi + \eta_2$$

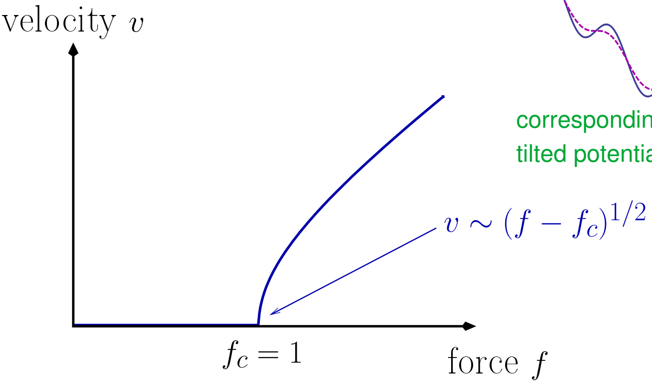
- Effective model: collective degrees of freedom

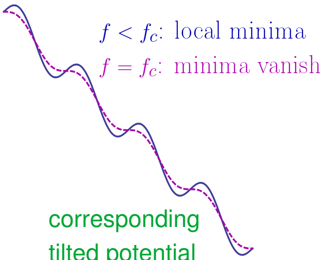
Position $r(t)$ coupled to phase $\phi(t)$

Depinning @ zero temperature

(1st case) Large $K_\perp$: ϕ decouples from r

$$\alpha\partial_t r = \underbrace{f - \cos\kappa r}_{\text{total force}}$$



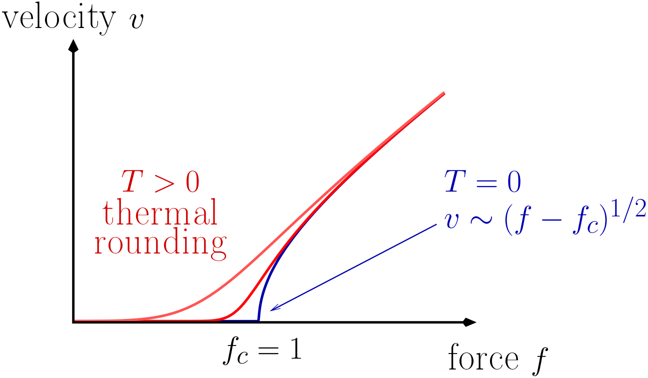


corresponding tilted potential

Depinning @ finite temperature

(1st case) Large $K_\perp$: ϕ decouples from r

$$\alpha\partial_t r = \underbrace{f - \cos\kappa r}_{\text{total force}} + \eta$$



Depinning @ zero temperature

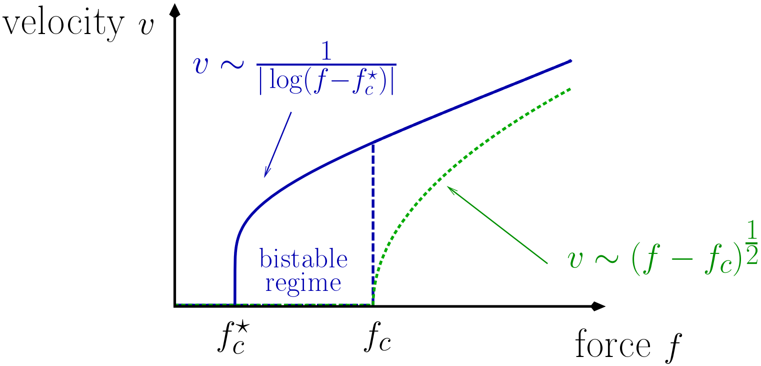
(2nd case) Small $K_\perp$: ϕ matters

$$\alpha\partial_t r - \partial_t\phi = f - \cos\kappa r$$
$$\alpha\partial_t\phi + \partial_t r = -\frac{1}{2}K_\perp \sin 2\phi$$

Depinning @ zero temperature

(2nd case) Small $K_\perp$: ϕ matters

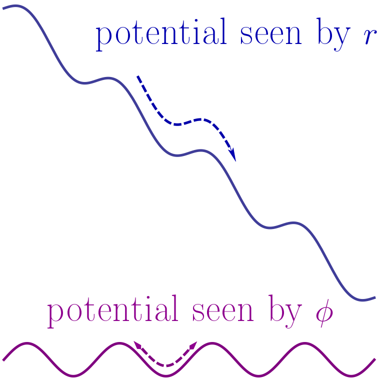
- Dramatic change in the depinning law: $v \sim \frac{1}{|\log(f - f_c^\star)|}$



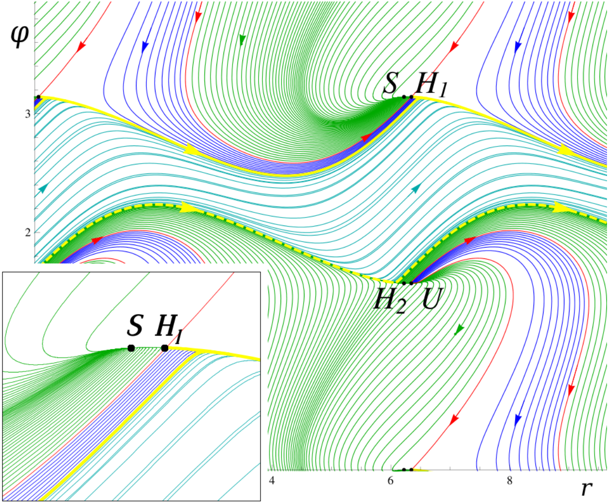
- Depinning at lower critical force: $f_c^\star < f_c$
- Bistability

Physical interpretation

In the bistable regime $f_c^\star < f < f_c$: ϕ helps r to cross barriers



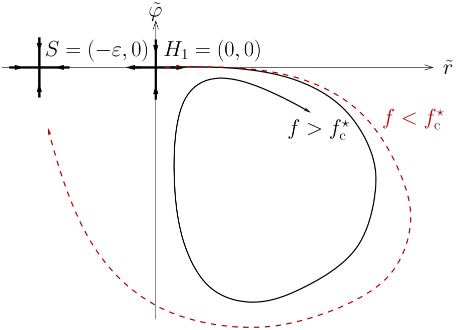
Phase space



In the bistable regime ($f_c^\star < f < f_c$)

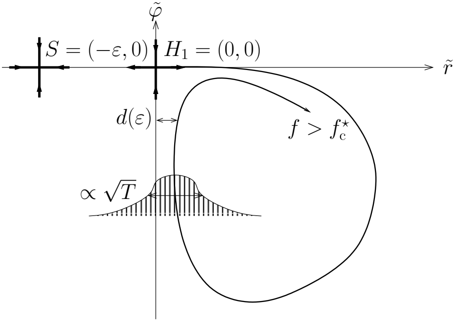
Phase space

Homoclinic bifurcation: $(\varepsilon \propto f_c - f > 0)$



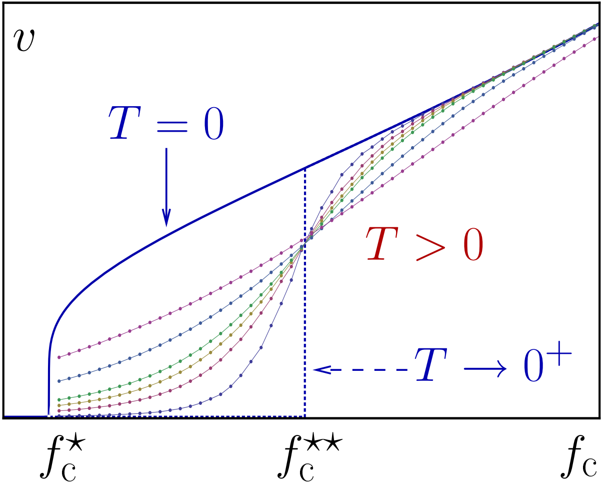
Phase space: $T > 0$

Homoclinic bifurcation with noise: $(\varepsilon \propto f_c - f > 0)$



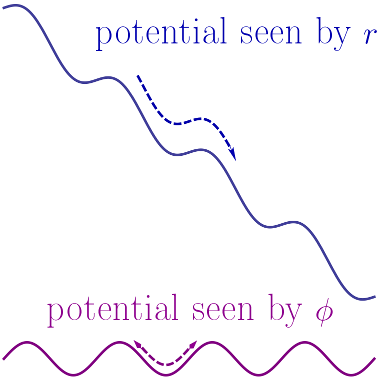
$$\text{escape time} \sim \underbrace{\exp\left(\frac{\varepsilon^3}{T}\right)}_{\text{Arrhenius}} \underbrace{\exp\left(-\frac{A}{T} d(\varepsilon)^2\right)}_{\text{Trapping probability}}$$

Finite temperature



Force-velocity characteristics

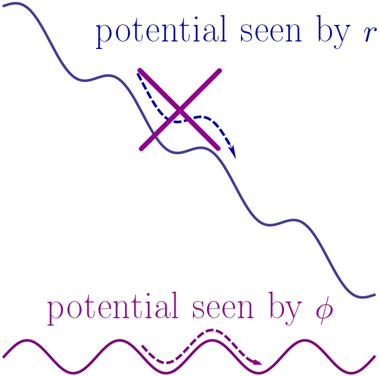
This is not the end of the story



The phase ϕ plays the role of inertia: helps to cross barriers

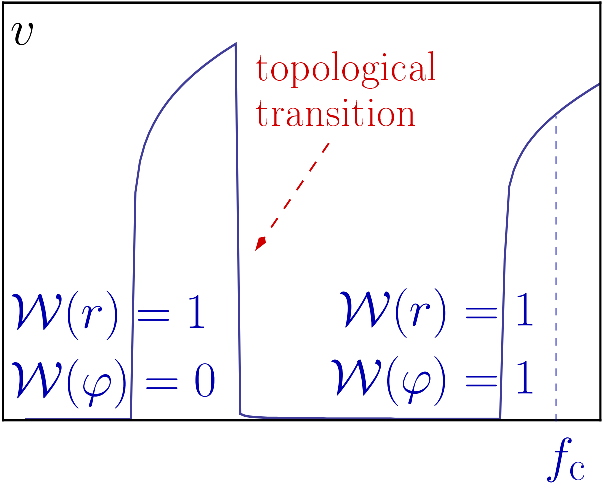
This is not the end of the story

(3rd case) Even smaller $K_\perp$



inertia is unbounded whereas ϕ is bounded and periodic

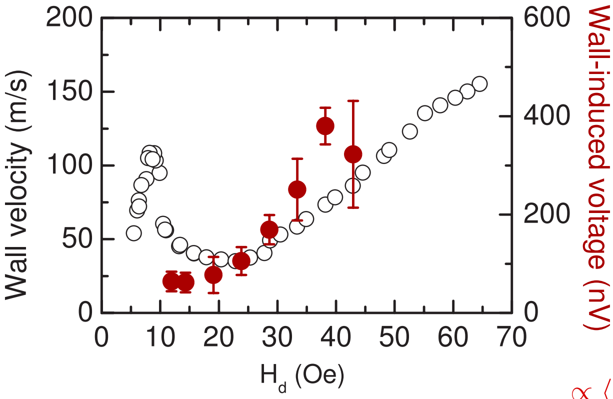
Topological transition



Successive regimes characterised by winding numbers $\mathcal{W}$

Experimental test? SPINTRONICS

experiment from Yang, Beach *et al.*, PRL **102** 067201 (2009)



$\propto \langle \partial_t \phi \rangle$

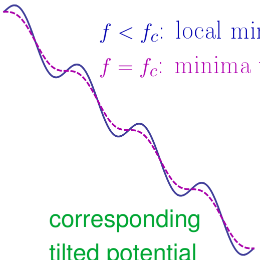
A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos \kappa r$$

($\Leftrightarrow$ Josephson junction)

A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos \kappa r$$

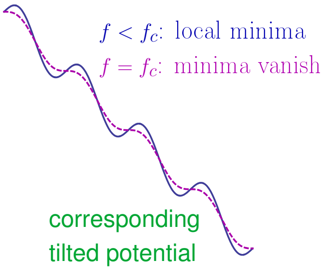


A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos\kappa r$$

Two regimes of mass:

- Small masses ($m < m_c$):
  same depinning transition as for $m = 0$
- Larger masses ($m > m_c$):

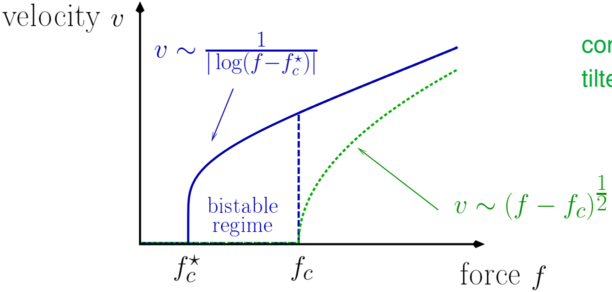


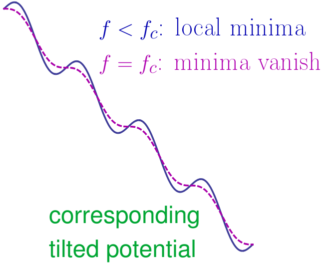
A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos\kappa r$$

Two regimes of mass:

- Small masses ($m < m_c$):
  same depinning transition as for $m = 0$
- Larger masses ($m > m_c$):



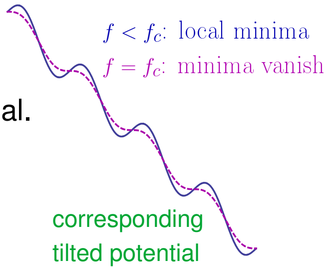


A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos \kappa r$$

Ongoing work [V. Purrello, AB. Kolton]:

- Criterion for the critical mass m_c?
  Dependency on the **details** of the potential.

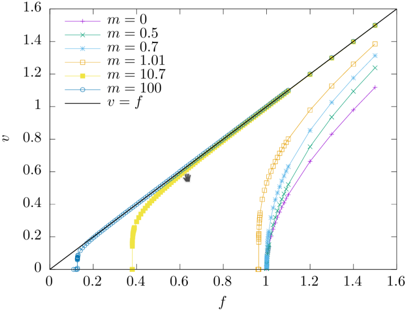


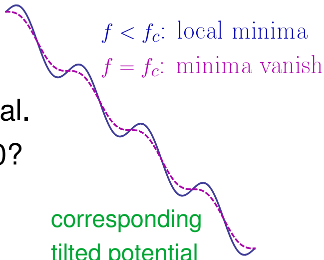
A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos\,\kappa r$$

Ongoing work [V. Purrello, AB. Kolton]:

- Criterion for the critical mass m_c?
  Dependency on the **details** of the potential.
- Scaling analysis in the regime $f_c - f_c^\star \to 0$?





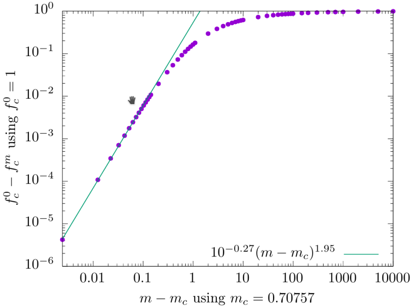
corresponding
tilted potential

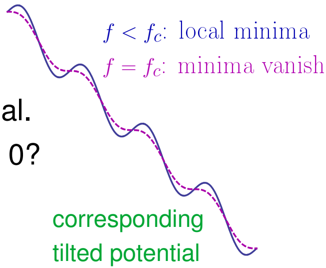
A correspondence: dynamics of a massive particle

$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos\,\kappa r$$

Ongoing work [V. Purrello, AB. Kolton]:

- Criterion for the critical mass m_c?
  Dependency on the **details** of the potential.
- Scaling analysis in the regime $m - m_c \to 0$?



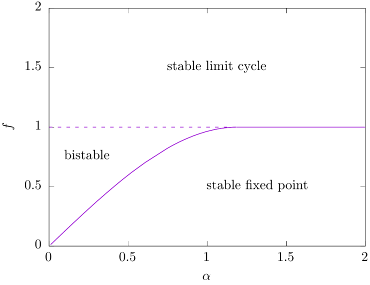


A correspondence: dynamics of a massive particle

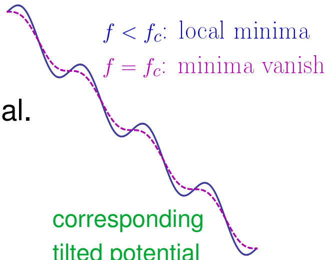
$$m\,\partial_t^2 r + \alpha\,\partial_t r = f - \cos\,\kappa r$$

Ongoing work [V. Purrello, AB. Kolton]:

- Criterion for the critical mass m_c?
  Dependency on the **details** of the potential.
- Phase diagram



$[\alpha = m^{-1/2}]$



corresponding tilted potential

Including elasticity (full 1D interface)

[Ongoing work with Alejandro Kolton, Victor Purrello]

$$m\,\partial_t^2 r + \alpha\,\partial_t r = c\partial_z^2 r + f - \partial_z V(r)$$

Interface with a **inertia** (velocity = internal degree of freedom)

- Is there a critical mass m_c? (i.e. $m > m_c \Rightarrow$ depinning changes)
  $\longrightarrow$ What is the phase diagram?
- How is the depinning affected?
  $\longrightarrow$ Bistability? Finite depinning exponent β?

Including elasticity (full 1D interface)

[Ongoing work with Alejandro Kolton, Victor Purrello]

$$m\,\partial_t^2 r + \alpha\,\partial_t r = c\partial_z^2 r + f - \partial_z V(r)$$

Interface with a **inertia** (velocity $=$ internal degree of freedom)

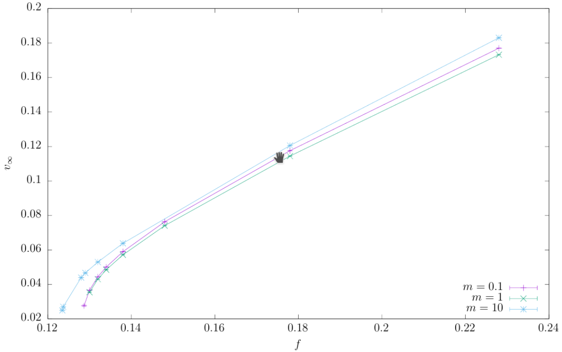
- Is there a critical mass m_c? (i.e. $m > m_c \Rightarrow$ depinning changes)
  $\longrightarrow$ What is the phase diagram?
- How is the depinning affected?
  $\longrightarrow$ Bistability? Finite depinning exponent β?
- How are geometrical properties affected?
- How, in presence of disorder, can there be a **cooperation** between position and velocity?
- What are the effects of temperature?

Including elasticity (full 1D interface)

[Ongoing work with Alejandro Kolton, Victor Purrello]

$$m\,\partial_t^2 r + \alpha\,\partial_t r = c\partial_z^2 r + f - \partial_z V(r)$$

Interface with a **inertia** (velocity = internal degree of freedom)



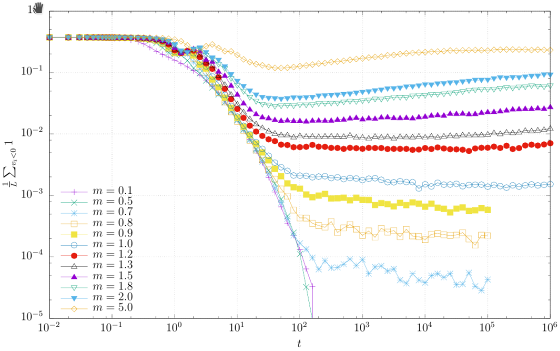
There exists a **bistable regime** ($f_c^\star < f_c$).

Including elasticity (full 1D interface)

[Ongoing work with Alejandro Kolton, Victor Purrello]

$$m\,\partial_t^2 r + \alpha\,\partial_t r = c\partial_z^2 r + f - \partial_z V(r)$$

Interface with a **inertia** (velocity = internal degree of freedom)



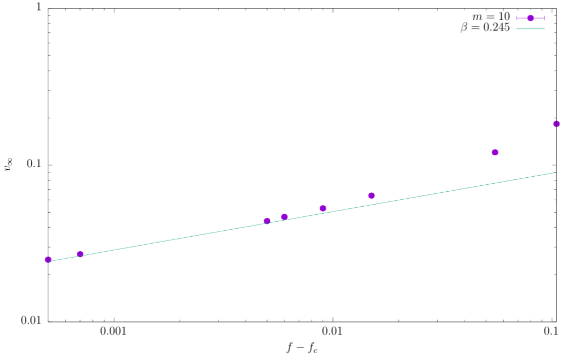
Cooperative motion $f_c^\star < f < f_c$.

Including elasticity (full 1D interface)

[Ongoing work with Alejandro Kolton, Victor Purrello]

$$m\,\partial_t^2 r + \alpha\,\partial_t r = c\partial_z^2 r + f - \partial_z V(r)$$

Interface with a **inertia** (velocity = internal degree of freedom)



Same depinning exponent for $m < m_c$ and $m = 0$.

Summary

Internal degree of freedom

- unusual depinning law
- bistability at zero T
- non-monotonous $v(f)$ at finite T

Perspective

- Interface with elasticity $\longleftrightarrow$ modified creep law?
- Current-driven wall $\longleftrightarrow$ spintronics?
- Experiments $\longleftrightarrow$ periodic patterning?
- Other internal degrees $\longleftrightarrow$ coupled interfaces?

Summary

Internal degree of freedom

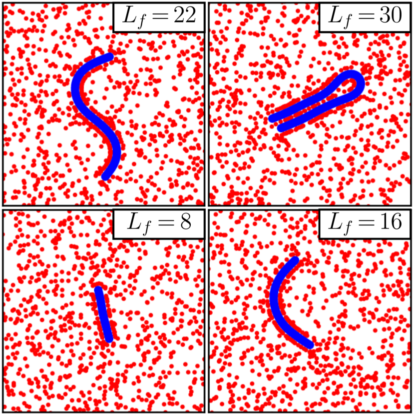
- unusual depinning law
- bistability at zero T
- non-monotonous $v(f)$ at finite T

Perspective

- Interface with elasticity $\longleftrightarrow$ modified creep law?
- Current-driven wall $\longleftrightarrow$ spintronics?
- Experiments $\longleftrightarrow$ periodic patterning?
- Other internal degrees $\longleftrightarrow$ coupled interfaces?

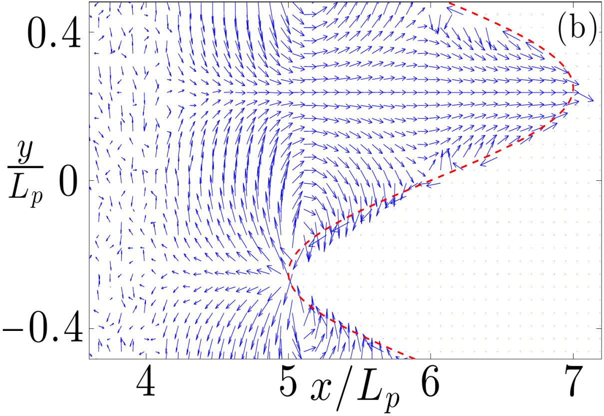
Interfaces in active materials

[Nikolai Nikola, Alexandre P. Solon, Yariv Kafri, Mehran Kardar, Julien Tailleur, Raphaël Voituriez
PRL **117** 098001 (2016)]



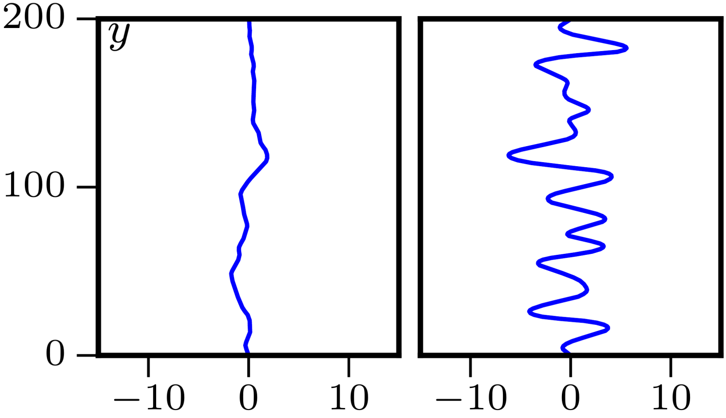
Interfaces in active materials

[Nikolai Nikola, Alexandre P. Solon, Yariv Kafri, Mehran Kardar, Julien Tailleur, Raphaël Voituriez
PRL **117** 098001 (2016)]



Interfaces in active materials

[Nikolai Nikola, Alexandre P. Solon, Yariv Kafri, Mehran Kardar, Julien Tailleur, Raphaël Voituriez
PRL **117** 098001 (2016)]



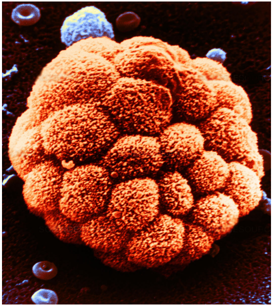
Interfaces in active materials

[Nikolai Nikola, Alexandre P. Solon, Yariv Kafri, Mehran Kardar, Julien Tailleur, Raphaël Voituriez
PRL **117** 098001 (2016)]

- **Issues**:
  - Energy is injected and dissipated in the bulk.
  - Usual concepts of thermodynamics (pressure,. . .) do not apply.

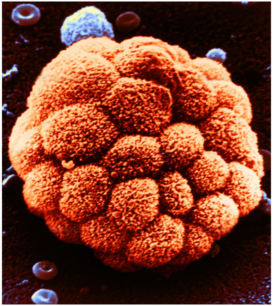
- **Questions**:
  - How to understand fluctuations of interfaces/walls?
  - How to build effective models?

Interfaces dynamics of tissues



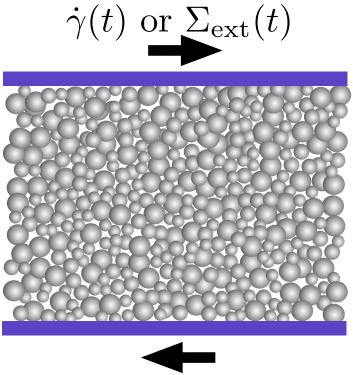
$\xrightarrow{?}$

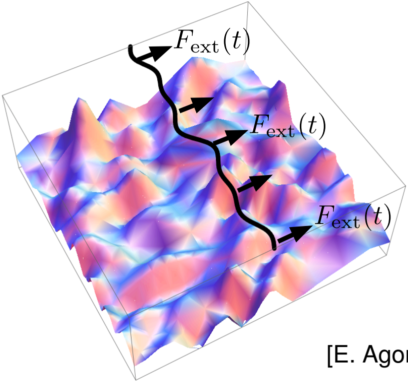
Interfaces dynamics of tissues



$\overset{?}{\longrightarrow}$

Relation between yielding and depinning

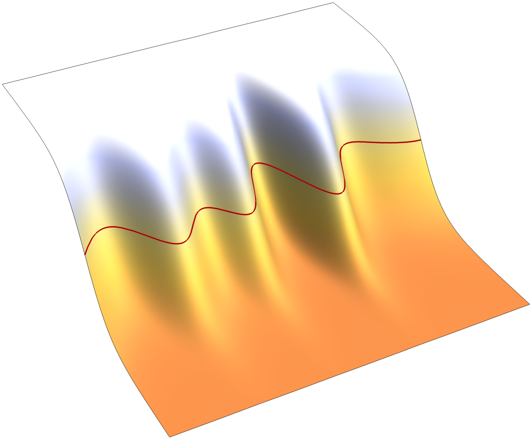




[E. Agoritsas]

- Yielding transition $\leftrightarrow$ depinning transition
- Burst dynamics $\leftrightarrow$ avalanche dynamics

Relation between bulk and effective descriptions



[Nirvana Belén Caballero, Elisabeth Agoritsas, Thierry Giamarchi]

Thank you for your attention!

Additional slides

Equilibrium directed polymer at temperature T (and $f = 0$)

- Elasticity: tends to flatten the interface

$$\mathcal{H}^{\text{el}}[y(t), t_{\text{f}}] = \frac{c}{2} \int_0^{t_{\text{f}}} dt\, \big[\partial_t y(t)\big]^2$$

- Disorder: tends to bend it

$$\mathcal{H}_V^{\text{dis}}[y(t), t_{\text{f}}] = \int_0^{t_{\text{f}}} dt\, V\big(t, y(t)\big)$$

Competition btw “order” and “disorder”

Equilibrium directed polymer at temperature T (and $f = 0$)

- Elasticity: tends to flatten the interface

$$\mathcal{H}^{\mathrm{el}}[y(t), t_{\mathrm{f}}] = \frac{c}{2} \int_0^{t_{\mathrm{f}}} dt \, \left[\partial_t y(t)\right]^2$$

- Disorder: tends to bend it

$$\mathcal{H}_V^{\mathrm{dis}}[y(t), t_{\mathrm{f}}] = \int_0^{t_{\mathrm{f}}} dt \, V\big(t, y(t)\big)$$

weight of $(y(t))_{0<t<t_{\mathrm{f}}}$:

$$e^{-\mathcal{H}_V / T}$$

with

$$\mathcal{H}_V = \mathcal{H}^{\mathrm{el}} + \mathcal{H}_V^{\mathrm{dis}}$$

$\boxed{\text{Competition btw “order” and “disorder”}}$

- Interpretations of the elastic Hamiltonian:

elasticity or kinetic energy or Wiener measure

Equilibrium directed polymer at temperature T (and $f=0$)

- Elasticity: tends to flatten the interface

$$\mathcal{H}^{\mathrm{el}}[y(t),t_{\mathrm{f}}] = \frac{c}{2}\int_0^{t_{\mathrm{f}}} dt\, \left[\partial_t y(t)\right]^2$$

- Disorder: tends to bend it

$$\mathcal{H}_V^{\mathrm{dis}}[y(t),t_{\mathrm{f}}] = \int_0^{t_{\mathrm{f}}} dt\, V(t,y(t))$$

weight of $(y(t))_{0<t<t_{\mathrm{f}}}$:

$$e^{-\mathcal{H}_V/T}$$

with

$$\mathcal{H}_V = \mathcal{H}^{\mathrm{el}} + \mathcal{H}_V^{\mathrm{dis}}$$

Competition btw "order" and "disorder"

- Ingredients up to now:

elastic constant c disorder potential $V(t,y)$ temperature T

A naive scaling argument 1/3

Roughness function (variance of the end-point fluctuations):

$$B(t_{\mathrm{f}}) = \overline{\langle y(t_{\mathrm{f}})^2 \rangle} \qquad \overline{\ \cdot\ } = \text{disorder average} \qquad \langle \cdot \rangle = \text{thermal average}$$

Path-integral writing:

$$B(t_{\mathrm{f}}) = \int \mathcal{D}V \ \mathbb{P}[V] \, \frac{\int_{y(0)=0} \mathcal{D}y(t)\, y(t_{\mathrm{f}})^2 \, \mathrm{e}^{-\frac{1}{T}\mathcal{H}[y(t),V;t_{\mathrm{f}}]}}{\int_{y(0)=0} \mathcal{D}y(t)\, \mathrm{e}^{-\frac{1}{T}\mathcal{H}[y(t),V;t_{\mathrm{f}}]}}$$

Uncorrelated disorder: Gaussian, centered, with

$$\overline{V(t,y)V(t',y')} = D\,\delta(t'-t)\delta(y'-y)$$

which rescales as

$$V(b\hat{t}, a\hat{y}) \stackrel{d}{=} a^{-\frac{1}{2}} b^{-\frac{1}{2}} D^{\frac{1}{2}} \hat{V}(\hat{t},\hat{y})$$

A naive scaling argument 2/3

Flory rescaling: $t = t_{\mathrm{f}}\,\hat{t}, \qquad y = t_{\mathrm{f}}^{\zeta_{\mathrm{F}}} \big(\frac{D}{c^2}\big)^{\frac{1}{5}}\hat{y}, \qquad \zeta_{\mathrm{F}} = \frac{3}{5}$
ensures

$$\frac{1}{T}\mathcal{H}^{\text{el}} = \frac{(cD^2)^{1/5}}{T} t_{\mathrm{f}}^{1/5}\, \frac{1}{2}\int_0^1 d\hat{t}\, \big[\partial_{\hat{t}}\hat{y}(\hat{t})\big]^2$$

$$\frac{1}{T}\mathcal{H}^{\text{dis}} \overset{d}{=} \frac{(cD^2)^{1/5}}{T} t_{\mathrm{f}}^{1/5} \int_0^1 d\hat{t}\, \hat{V}\big(\hat{t}, \hat{y}(\hat{t})\big)$$

with correlations:

$$\overline{\hat{V}(\hat{t},\hat{y})\hat{V}(\hat{t}',\hat{y}')} = \delta(\hat{t}'-\hat{t})\delta(\hat{y}'-\hat{y})$$

A naive scaling argument 2/3

Flory rescaling: $t = t_{\rm f}\,\hat t,\qquad y = t_{\rm f}^{\zeta_{\rm F}}\left(\frac{D}{c^2}\right)^{\frac15}\hat y,\qquad \zeta_{\rm F}=\frac35$

Rescaling of the roughness

$$B(t_{\rm f}) \;=\; \left[\frac{D}{c^2}\right]^{2/5} t_{\rm f}^{6/5}\,\hat B(t_{\rm f})$$

$$\hat B(t_{\rm f}) \;=\; \frac{\displaystyle\int_{\hat y(0)=0}\mathcal{D}\hat y(\hat t)\;\hat y(1)^2\exp\Big\{-\frac{(cD^2)^{\frac15}}{T}t_{\rm f}^{\frac15}\int_0^1 d\hat t\,\Big[\tfrac12(\partial_{\hat t}\hat y)^2+\hat V(\hat t,\hat y(\hat t))\Big]\Big\}}{\displaystyle\int_{\hat y(0)=0}\mathcal{D}\hat y(\hat t)\exp\Big\{-\frac{(cD^2)^{\frac15}}{T}t_{\rm f}^{\frac15}\int_0^1 d\hat t\,\Big[\tfrac12(\partial_{\hat t}\hat y)^2+\hat V(\hat t,\hat y(\hat t))\Big]\Big\}}$$

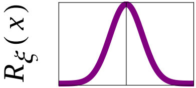
If saddle-point trajectory $\hat y^\star(\hat t)$ exists at $t_{\rm f}\to\infty$, **it is indep. of $t_{\rm f}$**

$$B(t_{\rm f}) \;=\; \left[\frac{D}{c^2}\right]^{2/5} t_{\rm f}^{6/5}\,\overline{\hat y^\star(1)^2}\qquad\text{not the expected KPZ behaviour}\quad\sim t_{\rm f}^{4/3}$$

A naive scaling argument 3/3

Where is t_f? In the disorder correlations on a lengthscale ξ:

$$\overline{V(z,x)V(z',x')} = D\,\delta(z'-z)R_\xi(x'-x)$$



scaling as $R_\xi(x) = \frac{1}{\xi} R_{\xi=1}(x/\xi)$

A naive scaling argument 3/3

Where is t_{f}? In the disorder correlations on a lengthscale ξ:

$$\overline{V(z,x)V(z',x')} = D\,\delta(z'-z)R_\xi(x'-x)$$

Flory rescaling: $t = t_{\mathrm{f}}\,\hat{t}, \qquad y = t_{\mathrm{f}}^{\zeta_{\mathrm{F}}}\left(\frac{D}{c^2}\right)^{\frac{1}{5}}\hat{y}, \qquad \zeta_{\mathrm{F}} = \frac{3}{5}$

Rescaling of the roughness

$$B(t_{\mathrm{f}}) = \left[\tfrac{D}{c^2}\right]^{2/5} t_{\mathrm{f}}^{6/5}\,\hat{B}(t_{\mathrm{f}}) \qquad \text{where denoting} \quad \hat{\xi}(t_{\mathrm{f}}) = \frac{\xi}{t_{\mathrm{f}}^{\zeta_{\mathrm{F}}}\left(\frac{D}{c^2}\right)^{\frac{1}{5}}}$$

$$\hat{B}(t_{\mathrm{f}}) = \frac{\displaystyle\int_{\hat{y}(0)=0}\mathcal{D}\hat{y}(\hat{t})\,\hat{y}(1)^2\exp\left\{-\frac{(cD^2)^{\frac{1}{5}}}{T}t_{\mathrm{f}}^{\frac{1}{5}}\int_0^1 d\hat{t}\left[\tfrac{1}{2}(\partial_{\hat{t}}\hat{y})^2+\hat{V}_{\hat{\xi}(t_{\mathrm{f}})}(\hat{t},\hat{y}(\hat{t}))\right]\right\}}{\displaystyle\int_{\hat{y}(0)=0}\mathcal{D}\hat{y}(\hat{t})\exp\left\{-\frac{(cD^2)^{\frac{1}{5}}}{T}t_{\mathrm{f}}^{\frac{1}{5}}\int_0^1 d\hat{t}\left[\tfrac{1}{2}(\partial_{\hat{t}}\hat{y})^2+\hat{V}_{\hat{\xi}(t_{\mathrm{f}})}(\hat{t},\hat{y}(\hat{t}))\right]\right\}}$$

$$B(t_{\mathrm{f}}) = \left[\tfrac{D}{c^2}\right]^{2/5} t_{\mathrm{f}}^{6/5}\,\overline{\hat{y}^\star(1)^2} \qquad \left\{\begin{array}{l}\text{the } t_{\mathrm{f}}\to\infty \text{ saddle-point trajectory}\\ \hat{y}^\star(\hat{t}) \textbf{ depends on } t_{\mathrm{f}} \textbf{ through } \hat{\xi}(t_{\mathrm{f}})\end{array}\right\}$$

A naive scaling argument 3/3

Where is t_{f}? In the disorder correlations on a lengthscale ξ:

$$\overline{V(z,x)V(z',x')} = D\,\delta(z'-z)R_\xi(x'-x)$$

Flory rescaling: $t = t_{\mathrm{f}}\,\hat{t}, \qquad y = t_{\mathrm{f}}^{\zeta_{\mathrm{F}}}\big(\frac{D}{c^2}\big)^{\frac{1}{5}}\hat{y}, \qquad \zeta_{\mathrm{F}} = \frac{3}{5}$

Rescaling of the roughness

$$B(t_{\mathrm{f}}) = \big[\tfrac{D}{c^2}\big]^{2/5} t_{\mathrm{f}}^{6/5}\,\hat{B}(t_{\mathrm{f}}) \qquad \text{where denoting} \qquad \hat{\xi}(t_{\mathrm{f}}) = \frac{\xi}{t_{\mathrm{f}}^{\zeta_{\mathrm{F}}}\big(\frac{D}{c^2}\big)^{\frac{1}{5}}}$$

$$\hat{B}(t_{\mathrm{f}}) = \frac{\overline{\int_{\hat{y}(0)=0}\mathcal{D}\hat{y}(\hat{t})\,\hat{y}(1)^2\exp\Big\{-\frac{(cD^2)^{\frac{1}{5}}}{T}t_{\mathrm{f}}^{\frac{1}{5}}\int_0^1 d\hat{t}\,\Big[\frac{1}{2}(\partial_{\hat{t}}\hat{y})^2+\hat{V}_{\hat{\xi}(t_{\mathrm{f}})}(\hat{t},\hat{y}(\hat{t}))\Big]\Big\}}}{\int_{\hat{y}(0)=0}\mathcal{D}\hat{y}(\hat{t})\exp\Big\{-\frac{(cD^2)^{\frac{1}{5}}}{T}t_{\mathrm{f}}^{\frac{1}{5}}\int_0^1 d\hat{t}\,\Big[\frac{1}{2}(\partial_{\hat{t}}\hat{y})^2+\hat{V}_{\hat{\xi}(t_{\mathrm{f}})}(\hat{t},\hat{y}(\hat{t}))\Big]\Big\}}$$

$$B(t_{\mathrm{f}}) = \big[\tfrac{D}{c^2}\big]^{2/5} t_{\mathrm{f}}^{6/5}\,\overline{\hat{y}^\star(1)^2} \qquad \Big\{\cdots\Big\} \Rightarrow \text{modifies the exponent } \frac{6}{5}\mapsto\frac{4}{3}$$

Free-energy fluctuations of trajectories $(0,0) \rightsquigarrow (t,y)$

How to solve this issue?

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

Free-energy fluctuations of trajectories $(0,0) \rightsquigarrow (t,y)$

How to solve this issue?

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

- **Stochastic Heat Equation** (Feynman-Kac formula)

$$\partial_t Z_V = \Big[\frac{T}{2c}\partial_y^2 - \frac{1}{T}V(t,y)\Big] Z_V(t,y) \qquad \textbf{(SHE)}$$

  Linear, multiplicative noise, reversible

- **Kardar-Parisi-Zhang** equation

$$\partial_t F_V = \frac{T}{2c}\partial_y^2 F_V - \frac{1}{2c}\big[\partial_y F_V\big]^2 + V(t,y) \qquad \textbf{(KPZ)}$$

  Non-linear, additive noise, non-reversible

  $F_V(t,y) \equiv$ interface height at position y and time t

Free-energy fluctuations of trajectories $(0,0) \rightsquigarrow (t,y)$

How to solve this issue?

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

- **Stochastic Heat Equation** (Feynman-Kac formula)

$$\partial_t Z_V = \Big[\frac{T}{2c}\partial_y^2 - \frac{1}{T}V(t,y)\Big]Z_V(t,y) \qquad \textbf{(SHE)}$$

Linear, multiplicative noise, $Z_V(0,y) = \delta(y)$

- **Kardar-Parisi-Zhang** equation

$$\partial_t F_V = \frac{T}{2c}\partial_y^2 F_V - \frac{1}{2c}\big[\partial_y F_V\big]^2 + V(t,y) \qquad \textbf{(KPZ)}$$

Non-linear, additive noise, $F_V(0,y)$: "sharp wedge" initial cond.

$F_V(t,y) \equiv$ interface height at position y and time t

Statistical tilt symmetry

How to solve this issue?

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

Statistical tilt symmetry

How to solve this issue?

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

- **Statistical Tilt Symmetry**

$$\boxed{F_V(t,y) = \underbrace{c\frac{y^2}{2t} + \frac{T}{2}\log\frac{2\pi Tt}{c}}_{\substack{\text{thermal contribution}\\ F_{V\equiv 0}}} + \underbrace{\bar{F}_V(t,y)}_{\substack{\text{disorder}\\ \text{contribution}}}} \qquad (\textbf{STS})$$

- **Tilted** KPZ equation for $\bar{F}_V(t,y)$

$$\partial_t\bar{F}_V + \frac{y}{t}\partial_y\bar{F}_V = \frac{T}{2c}\partial_y^2\bar{F}_V - \frac{1}{2c}\left[\partial_y\bar{F}_V\right]^2 + V(t,y)$$

Non-linear, additive noise, $\bar{F}_V(0,y)\equiv 0$: “simple” initial cond.

Known results $@\xi = 0$ $[\Longleftrightarrow T \to \infty \; @\xi > 0]$

- **Infinite-time limit** $t_\mathrm{f} \to \infty$ (steady state)

  $\bar{F}(t_\mathrm{f} = \infty, y)$ distributed as a 2-sided Brownian Motion

  i.e.: $\mathbb{P}\big[\bar{F}(t_\mathrm{f} = \infty, y)\big]$ Gaussian, of correlator

$$\overline{\big[\bar{F}(t_\mathrm{f} = \infty, y) - \bar{F}(t_\mathrm{f} = \infty, y')\big]^2} = \widetilde{D}\,|y - y'| \quad \text{with} \quad \boxed{\widetilde{D} = \frac{cD}{T}}$$

Known results @$\xi = 0$ $\quad [\Longleftrightarrow T \to \infty$ @$\xi > 0]$

- **Infinite-time limit** $t_\mathrm{f} \to \infty$ (steady state)

  $\bar{F}(t_\mathrm{f} = \infty, y)$ distributed as a 2-sided Brownian Motion

  i.e.: $\mathbb{P}\big[\bar{F}(t_\mathrm{f} = \infty, y)\big]$ Gaussian, of correlator

$$\overline{\big[\bar{F}(t_\mathrm{f} = \infty, y) - \bar{F}(t_\mathrm{f} = \infty, y')\big]^2} = \widetilde{D}\,|y - y'| \quad \text{with} \quad \boxed{\widetilde{D} = \frac{cD}{T}}$$

- Rescaling of the disorder free-energy $\bar{F}(t_\mathrm{f} = \infty, a\hat{y})$

$$\bar{F}(t_\mathrm{f} = \infty, a\hat{y}) \stackrel{d}{=} a^{1/2}\widetilde{D}^{1/2} \underbrace{\hat{F}(t_\mathrm{f} = \infty, \hat{y})}_{\widetilde{D}=1}$$

Saddle-point argument for the KPZ exponent

Flory rescaling for the free-energy

$$t = t_{\mathrm{f}}\hat{t}, \qquad y = (\widetilde{D}/c^2)^{\frac{1}{3}}\, t_{\mathrm{f}}^{\frac{2}{3}}\, \hat{y}, \qquad \bar{F}_V(t,y) \overset{(d)}{=} (\widetilde{D}^2 t_{\mathrm{f}}/c)^{1/3} \hat{F}(\hat{t},\hat{y})$$

Rescaling of the roughness

$$B(t_{\mathrm{f}}) \underset{t_{\mathrm{f}}\to\infty}{\sim} \Big[\frac{\widetilde{D}}{c^2}\Big]^{\frac{2}{3}} t_{\mathrm{f}}^{\frac{4}{3}}\ \hat{B}(t_{\mathrm{f}})$$

$$\hat{B}(t_{\mathrm{f}}) = \overline{\frac{\displaystyle\int_{\mathbb{R}} d\hat{y}\ \hat{y}^2 \exp\Big\{-\tfrac{1}{T}\big(\tfrac{\widetilde{D}^2}{c}t_{\mathrm{f}}\big)^{\frac{1}{3}}\Big[\tfrac{\hat{y}^2}{2}+\hat{F}(\hat{t},\hat{y})\Big]\Big\}}{\displaystyle\int_{\mathbb{R}} d\hat{y}\ \exp\Big\{-\tfrac{1}{T}\big(\tfrac{\widetilde{D}^2}{c}t_{\mathrm{f}}\big)^{\frac{1}{3}}\Big[\tfrac{\hat{y}^2}{2}+\hat{F}(\hat{t},\hat{y})\Big]\Big\}}}$$

The $t_{\mathrm{f}} \to \infty$ saddle-point $[\hat{B}(t_{\mathrm{f}}) \sim \overline{(\hat{y}^\star)^2} \sim t_{\mathrm{f}}^0]$ gives the correct exponent

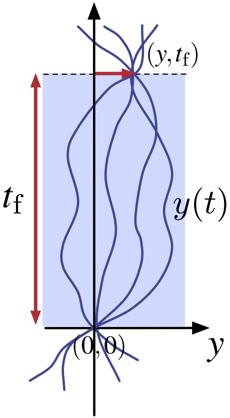
A simple picture of creep?

1D interface in the Directed Polymer (DP) language

- No bubbles
- No overhangs
- Interface lengthscale z

$\updownarrow$

DP 'time' t_{f}



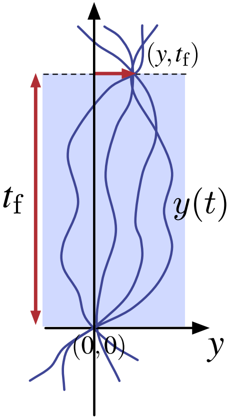
working at fixed 'time' t_{f} $\iff$
integration of fluctuations at scales smaller than t_{f}

1D interface in the Directed Polymer (DP) language

- No bubbles
- No overhangs
- Interface lengthscale z

$\updownarrow$

DP 'time' t_{f}



working at fixed 'time' t_{f} $\Longleftrightarrow$
integration of fluctuations at scales smaller than t_{f}

$$\boxed{\text{lengthscale} \equiv \text{time duration}}$$

Free-energy fluctuations of trajectories $(0,0) \rightsquigarrow (t,y)$

Time evolution as "renormalisation"

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

Free-energy fluctuations of trajectories $(0,0) \rightsquigarrow (t,y)$

Time evolution as "renormalisation"

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

- **Stochastic Heat Equation** (Feynman-Kac formula)

$$\partial_t Z_V = \Big[\frac{T}{2c}\partial_y^2 - \frac{1}{T}V(t,y)\Big] Z_V(t,y) \qquad \textbf{(SHE)}$$

Linear, multiplicative noise, reversible

- **Kardar-Parisi-Zhang** equation

$$\partial_t F_V = \frac{T}{2c}\partial_y^2 F_V - \frac{1}{2c}\big[\partial_y F_V\big]^2 + V(t,y) \qquad \textbf{(KPZ)}$$

Non-linear, additive noise, non-reversible

Free-energy fluctuations of trajectories $(0,0) \rightsquigarrow (t,y)$

Time evolution as “renormalisation”

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

- **Stochastic Heat Equation** (Feynman-Kac formula)

$$\partial_t Z_V = \Big[\frac{T}{2c}\partial_y^2 - \frac{1}{T}V(t,y)\Big]Z_V(t,y) \qquad \textbf{(SHE)}$$

  Linear, multiplicative noise

- **Kardar-Parisi-Zhang** equation

$$\partial_t F_V = \frac{T}{2c}\partial_y^2 F_V - \frac{1}{2c}\big[\partial_y F_V\big]^2 + V(t,y) \qquad \textbf{(KPZ)}$$

  Non-linear, additive noise
  [$F_V(t,y) \equiv$ interface height at position y and time t]

Statistical tilt symmetry

Time evolution as "renormalisation"

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t')\mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

Statistical tilt symmetry

Time evolution as "renormalisation"

- Partition function Z_V *vs.* Free-energy F_V

$$Z_V(t,y) = \int_{y(0)=0}^{y(t)=y} \mathcal{D}y(t') \mathrm{e}^{-\frac{1}{T}\mathcal{H}_V[y(t'),t]} \qquad F_V(t,y) = -\frac{1}{T}\log Z_V(t,y)$$

- **Statistical Tilt Symmetry**

$$\boxed{F_V(t,y) = \underbrace{c\frac{y^2}{2t} + \frac{T}{2}\log\frac{2\pi Tt}{c}}_{\substack{\text{thermal contribution}\\ F_{V\equiv 0}}} + \underbrace{\bar{F}_V(t,y)}_{\substack{\text{disorder}\\ \text{contribution}}}} \qquad \textbf{(STS)}$$

Known result

- **Infinite-time limit** $t_\mathrm{f} \to \infty$ (steady state)

  $\bar{F}(t_\mathrm{f} = \infty, y)$ distributed as a Brownian motion along y

  i.e. : $\mathbb{P}\big[\bar{F}(t_\mathrm{f} = \infty, y)\big]$ Gaussian, of correlator

$$\overline{\big[\bar{F}(t_\mathrm{f} = \infty, y) - \bar{F}(t_\mathrm{f} = \infty, y')\big]^2} = \widetilde{D}\,|y - y'|$$

Known result

- **Infinite-time limit** $t_{\mathrm{f}} \to \infty$ (steady state)

  $\bar{F}(t_{\mathrm{f}} = \infty, y)$ distributed as a Brownian motion along y

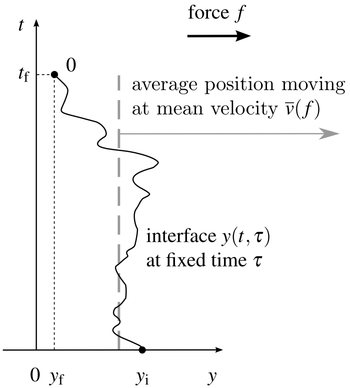
  i.e.: $\mathbb{P}\big[\bar{F}(t_{\mathrm{f}} = \infty, y)\big]$ Gaussian, of correlator

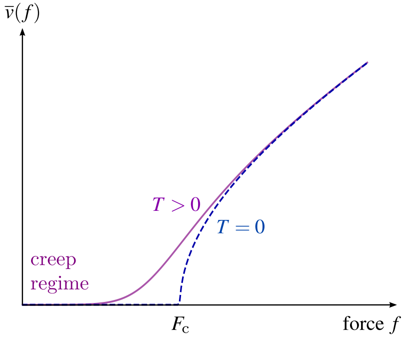
$$\overline{\big[\bar{F}(t_{\mathrm{f}} = \infty, y) - \bar{F}(t_{\mathrm{f}} = \infty, y')\big]^2} = \widetilde{D}\,|y - y'|$$

- Rescaling of the disorder free-energy $\bar{F}(t_{\mathrm{f}} = \infty, a\hat{y})$

$$\bar{F}(t_{\mathrm{f}} = \infty, a\hat{y}) \stackrel{d}{=} a^{1/2}\widetilde{D}^{1/2} \underbrace{\hat{F}(t_{\mathrm{f}} = \infty, \hat{y})}_{\widetilde{D}=1}$$

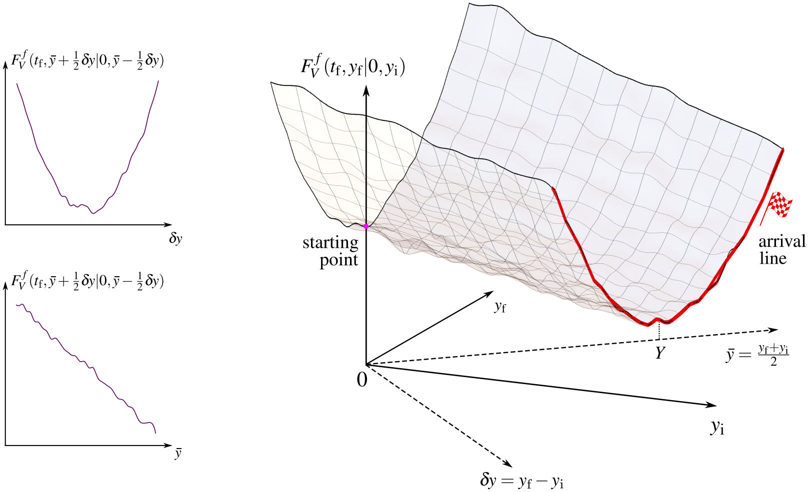
Non-linear response at small force





[Elisabeth Agoritsas, Reinaldo García-García, VL, Lev Truskinovsky and Damien Vandembroucq, J. Stat. Phys. **164** 1394 (2016)]

Effective model



Mean velocity $\longleftrightarrow$ Mean First Passage Time problem (MFPT)

Effective model

Mean velocity $\longleftrightarrow$ Mean First Passage Time problem (MFPT)

- Effective model at fixed t_f: quasistatic dynamics
  - ⋆ motion of a segment of length t_f
  - ⋆ extremities y_i and y_f follow Langevin dynamics
  - ⋆ forces derive from $F_V^f(t_\mathrm{f}, y_\mathrm{f} | t_\mathrm{i}, y_\mathrm{i})$
  - ⋆ exact at $f = 0$
- Optimisation over t_f at fixed f
  - ⋆ optimal t_f yielding the avalanche size at fixed f
  - ⋆ saddle-point argument after rescaling
  - ⋆ yields the creep law

$$\text{velocity} \sim \exp\Big\{ -\Big[\frac{\text{critical force}}{\text{force}}\Big]^{1/4} \Big\}$$

  - ⋆ creep exponent $\frac{1}{4}$ related to the KPZ exponent $\frac{2}{3}$