

Large deviations in (non-)equilibrium 1d transport models

Cécile Appert-Rolland¹, Thierry Bodineau², Bernard Derrida³,
Alberto Imparato⁴, Vivien Lecomte⁵, Frédéric van Wijland⁶

Julien Tailleur⁷, Jorge Kurchan⁸

¹LPT, Orsay ²DMA, Paris ³LPS, Paris ⁴DPA, Aarhus ⁵LPMA, Paris & DPMC, Genève
⁶MSC, Paris ⁷School of Physics, Edinburgh ⁸ESPCI, Paris

LPT&LPTMS – 30th September 2010



MaNEP
SWITZERLAND



**UNIVERSITÉ
DE GENÈVE**



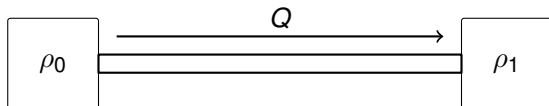
FNSNF

FONDS NATIONAL SUISSE
SCHWEIZERISCHER NATIONALFONDS
FONDO NAZIONALE SVIZZERO
SWISS NATIONAL SCIENCE FOUNDATION

Outline

- 1 Motivations
- 2 Microscopic approach
 - Operator approach
 - Bethe Ansatz
- 3 Macroscopic approach
 - Fluctuating hydrodynamics
 - Finite-size corrections
- 4 Mapping non-equilibrium to equilibrium
 - For the integrated current
 - For the density profile

Motivations



$$\text{Prob}(\mathcal{C}) \propto \exp \left\{ - \frac{\text{energy}(\mathcal{C})}{\text{temperature}} \right\} \quad \text{cannot describe}$$

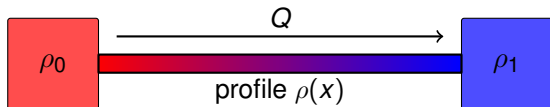
- Non-equilibrium steady-state

$$\text{Prob}[\rho(x)]$$

- Equilibrium fluctuations of dynamical observables

$$\text{Prob}[Q]$$

Motivations



$$\text{Prob}(\mathcal{C}) \propto \exp \left\{ - \frac{\text{energy}(\mathcal{C})}{\text{temperature}} \right\} \quad \text{cannot describe}$$

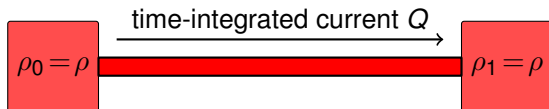
- Non-equilibrium steady-state

$$\text{Prob}[\rho(x)]$$

- Equilibrium fluctuations of dynamical observables

$$\text{Prob}[Q]$$

Motivations



$$\text{Prob}(\mathcal{C}) \propto \exp \left\{ - \frac{\text{energy}(\mathcal{C})}{\text{temperature}} \right\} \quad \text{cannot describe}$$

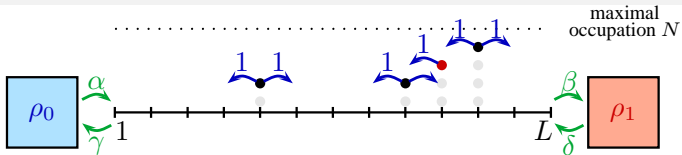
- Non-equilibrium steady-state

$$\text{Prob}[\rho(x)]$$

- Equilibrium fluctuations of dynamical observables

$$\text{Prob}[Q]$$

Exclusion Processes



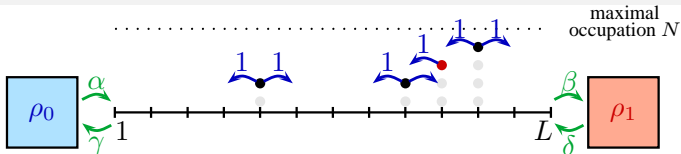
- Configurations: occupation numbers $\{n_i\}$
- Exclusion rule: $0 \leq n_i \leq N$
- Markov evolution

$$\partial_t P(\{n_i\}) = \sum_{n'_i} [W(n'_i \rightarrow n_i) P(\{n'_i\}) - W(n_i \rightarrow n'_i) P(\{n_i\})]$$

- Large deviation function of the time-integrated current Q

$$\langle e^{-sQ} \rangle \sim e^{t\psi(s)} \quad (\Leftrightarrow \text{determining } P(Q))$$

Exclusion Processes



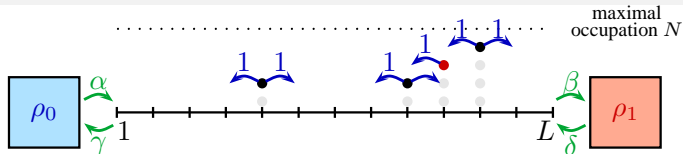
- Configurations: occupation numbers $\{n_i\}$
- Exclusion rule: $0 \leq n_i \leq N$
- Markov evolution

$$\partial_t P(\{n_i\}) = \sum_{n'_i} [W(n'_i \rightarrow n_i) P(\{n'_i\}) - W(n_i \rightarrow n'_i) P(\{n_i\})]$$

- Large deviation function of the time-integrated current Q

$$\langle e^{-sQ} \rangle \sim e^{t\psi(s)} \quad (\Leftrightarrow \text{determining } P(Q))$$

Exclusion Processes



- Configurations: occupation numbers $\{n_i\}$
- Exclusion rule: $0 \leq n_i \leq N$
- Markov evolution

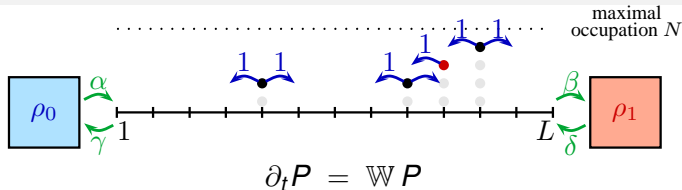
$$\partial_t P(\{n_i\}) = \sum_{n'_i} [W(n'_i \rightarrow n_i) P(\{n'_i\}) - W(n_i \rightarrow n'_i) P(\{n_i\})]$$

- Large deviation function of the time-integrated current Q

$$\langle e^{-sQ} \rangle \sim e^{t\psi(s)} \quad (\Leftrightarrow \text{determining } P(Q))$$

Operator representation

[Schütz & Sandow PRE 49 2726]



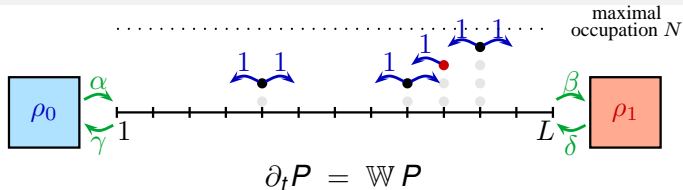
$$\mathbb{W} = \sum_{1 \leq k \leq L-1} [S_k^+ S_{k+1}^- + S_k^- S_{k+1}^+ - \hat{n}_k (1 - \hat{n}_{k+1}) - \hat{n}_{k+1} (1 - \hat{n}_k)] \\ + \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1] \\ + \delta [S_L^+ - (1 - \hat{n}_L)] + \beta [S_L^- - \hat{n}_L]$$

S^\pm and **creation and annihilation operators:**

$$S^+ |n\rangle = (N - n) |n + 1\rangle \quad S^- |n\rangle = n |n - 1\rangle$$

Operator representation

[Schütz & Sandow PRE 49 2726]



$$\mathbb{W} = \sum_{1 \leq k \leq L-1} [S_k^+ S_{k+1}^- + S_k^- S_{k+1}^+ - \hat{n}_k (1 - \hat{n}_{k+1}) - \hat{n}_{k+1} (1 - \hat{n}_k)]$$

$$+ \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1]$$

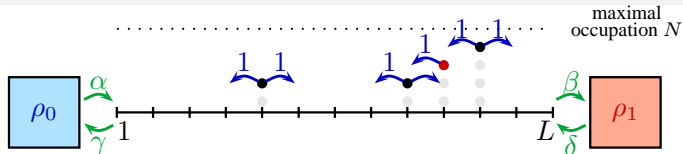
$$+ \delta [S_L^+ - (1 - \hat{n}_L)] + \beta [S_L^- - \hat{n}_L]$$

S^\pm and **creation and annihilation operators:**

$$S^+ |n\rangle = (N - n) |n + 1\rangle \quad S^- |n\rangle = n |n - 1\rangle$$

S^\pm and $S^z = \hat{n} - \frac{N}{2}$ are **spin operators** (with $j = \frac{N}{2}$)

Operator representation



$$\langle e^{-sQ} \rangle \sim e^{t\psi(s)} \quad \text{with} \quad \psi(s) = \max \text{Sp } \mathbb{W}(s)$$

$$\begin{aligned} \mathbb{W}(s) = & \sum_{1 \leq k \leq L-1} [S_k^+ S_{k+1}^- + S_k^- S_{k+1}^+ - \hat{n}_k(1 - \hat{n}_{k+1}) - \hat{n}_{k+1}(1 - \hat{n}_k)] \\ & + \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1] \\ & + \delta [S_L^+ e^s - (1 - \hat{n}_L)] + \beta [S_L^- e^{-s} - \hat{n}_L] \end{aligned}$$

Bethe Ansatz method [\[Appert, Derrida, VL, van Wijland, PRE 78 021122\]](#)

SSEP: maximal occupation $N = 1$
Periodic boundary conditions

Bethe Ansatz method [Appert, Derrida, VL, van Wijland, PRE 78 021122]

SSEP: maximal occupation $N = 1$

Periodic boundary conditions

Bethe Ansatz:

- eigenvector of components

$$\sum_{\mathcal{P}} \mathcal{A}(\mathcal{P}) \prod_{i=1}^{\mathcal{N}} [\zeta_{\mathcal{P}(i)}]^{x_i}$$

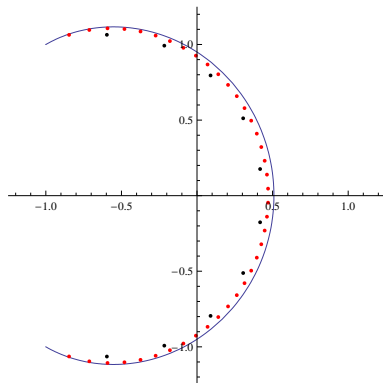
- eigenvalue

$$\psi(\mathbf{s}) = -2\mathcal{N} + \mathbf{e}^{-\mathbf{s}} [\zeta_1 + \dots + \zeta_{\mathcal{N}}] - \mathbf{e}^{\mathbf{s}} \left[\frac{1}{\zeta_1} + \dots + \frac{1}{\zeta_{\mathcal{N}}} \right]$$

- Bethe equations

$$\zeta_i^L = \prod_{\substack{j=1 \\ j \neq i}}^{\mathcal{N}} \left[-\frac{1 - 2\mathbf{e}^{-\mathbf{s}} \zeta_i + \mathbf{e}^{-2\mathbf{s}} \zeta_i \zeta_j}{1 - 2\mathbf{e}^{-\mathbf{s}} \zeta_j + \mathbf{e}^{-2\mathbf{s}} \zeta_i \zeta_j} \right]$$

Bethe Ansatz method [Appert, Derrida, VL, van Wijland, PRE **78** 021122]



Repartition of Bethe roots in the complex plane

Finite-size effects

[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

- large deviation function

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L}\rho(1-\rho)\mathbf{s}^2}_{\text{order 0}} + \underbrace{L^{-2}\mathcal{F}(u)}_{\text{finite-size}} \quad \text{with} \quad u = -\frac{1}{2}\rho(1-\rho)\mathbf{s}^2$$

- universal function

$$\mathcal{F}(u) = \sum_{k \geq 2} \frac{(-2u)^k B_{2k-2}}{\Gamma(k)\Gamma(k+1)}$$

- scaling of cumulants for the total current

$$\begin{aligned} \frac{1}{t} \langle Q^2 \rangle &\sim L \\ \frac{1}{t} \langle Q^{2k} \rangle &\sim L^{2k-2} \quad (k \geq 2) \end{aligned}$$

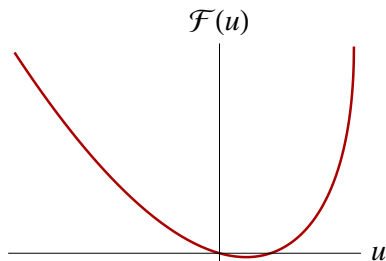
Finite-size effects

[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

- large deviation function

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L}\rho(1-\rho)\mathbf{s}^2}_{\text{order 0}} + \underbrace{L^{-2}\mathcal{F}(u)}_{\text{finite-size}} \quad \text{with} \quad u = -\frac{1}{2}\rho(1-\rho)\mathbf{s}^2$$

- universal function



Finite-size effects

[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

- large deviation function

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L}\rho(1-\rho)\mathbf{s}^2}_{\text{order 0}} + \underbrace{L^{-2}\mathcal{F}(u)}_{\text{finite-size}} \quad \text{with} \quad u = -\frac{1}{2}\rho(1-\rho)\mathbf{s}^2$$

- universal function

- technical remark:

. in that context (*symmetric* jump rates), the results of Kim [PRE **52** 3512] do not apply

. in particular the large L asymptotics reads

$$\frac{1}{L}\psi(L\mathbf{s}) = \frac{1}{2}\sigma\mathbf{s}^2 + \sigma^{3/2}|\mathbf{s}|^3$$

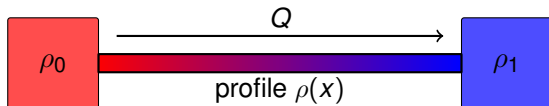
Outline

- 1 Motivations
- 2 Microscopic approach
 - Operator approach
 - Bethe Ansatz
- 3 Macroscopic approach
 - Fluctuating hydrodynamics
 - Finite-size corrections
- 4 Mapping non-equilibrium to equilibrium
 - For the integrated current
 - For the density profile



Macroscopic limit

[Tailleur, Kurchan, VL, JPA **41** 505001]

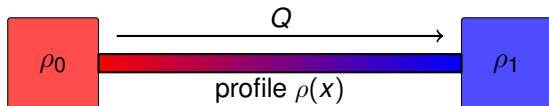


A reminder: propagator in quantum mechanics

$$\langle \text{final} | e^{it\mathbb{H}} | \text{initial} \rangle$$

Macroscopic limit

[Tailleur, Kurchan, VL, JPA **41** 505001]

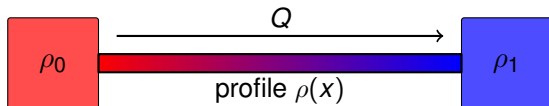


A reminder: propagator in quantum mechanics

$$\langle \text{final} | e^{it\mathbb{H}} | \text{initial} \rangle = \int dz_1 \dots dz_n \langle \text{final} | e^{i\Delta t\mathbb{H}} | \underline{z_n} \rangle \langle \underline{z_{n-1}} | e^{i\Delta t\mathbb{H}} | \underline{z_{n-2}} \rangle \dots \\ \dots \langle \underline{z_1} | e^{i\Delta t\mathbb{H}} | \text{initial} \rangle$$

Macroscopic limit

[Tailleur, Kurchan, VL, JPA 41 505001]

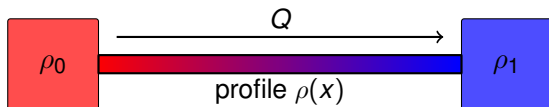


A reminder: propagator in quantum mechanics

$$\begin{aligned}
 \langle \text{final} | e^{it\mathbb{H}} | \text{initial} \rangle &= \int dz_1 \dots dz_n \langle \text{final} | e^{i\Delta t\mathbb{H}} | \underline{z}_n \rangle \langle \underline{z}_{n-1} | e^{i\Delta t\mathbb{H}} | \underline{z}_{n-2} \rangle \dots \\
 &\quad \dots \langle \underline{z}_1 | e^{i\Delta t\mathbb{H}} | \text{initial} \rangle \\
 &= \int \mathcal{D}p \mathcal{D}q \exp \left\{ i \frac{1}{\hbar} \underbrace{\mathcal{S}[p, q]}_{\text{action}} \right\}
 \end{aligned}$$

Macroscopic limit

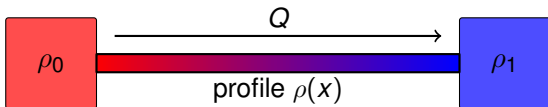
[Tailleur, Kurchan, VL, JPA 41 505001]

Using $SU(2)$ coherent states:

$$\langle \rho_f | e^{t\mathbb{W}} | \rho_i \rangle = \int_{\rho(0)=\rho_i}^{\rho(t)=\rho_f} \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L \underbrace{\mathcal{S}[\hat{\rho}, \rho]}_{\text{action}}\}$$

Macroscopic limit

[Tailleur, Kurchan, VL, JPA 41 505001]

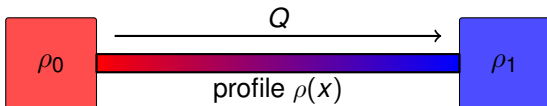
Using $SU(2)$ coherent states:

$$\langle \rho_f | e^{t\mathbb{W}} | \rho_i \rangle = \int_{\rho(0)=\rho_i}^{\rho(t)=\rho_f} \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L \underbrace{\mathcal{S}[\hat{\rho}, \rho]}_{\text{action}}\}$$

$$\langle e^{-sQ} \rangle \sim \langle \rho_f | e^{t\mathbb{W}_s} | \rho_i \rangle = \int_{\rho(0)=\rho_i}^{\rho(t)=\rho_f} \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L \underbrace{\mathcal{S}_s[\hat{\rho}, \rho]}_{\text{action}}\}$$

Macroscopic limit

[Tailleur, Kurchan, VL, JPA 41 505001]



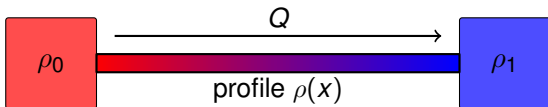
Same $\mathcal{S}_s[\hat{\rho}, \rho]$ as the MSR action of the Langevin evolution:

$$\partial_t \rho = -\partial_x [-\partial_x \rho + \xi]$$

$$\langle \xi(x, t) \xi(x', t') \rangle = \underbrace{\frac{1}{L} \rho(1 - \rho)}_{\text{density-dependent}} \delta(x' - x) \delta(t' - t)$$

Macroscopic limit

[Tailleur, Kurchan, VL, JPA 41 505001]



Same $\mathcal{S}_s[\hat{\rho}, \rho]$ as the MSR action of the Langevin evolution:

$$\partial_t \rho = -\partial_x [-\partial_x \rho + \xi]$$

$$\langle \xi(x, t) \xi(x', t') \rangle = \underbrace{\frac{1}{L} \rho(1 - \rho)}_{\text{density-dependent}} \delta(x' - x) \delta(t' - t)$$

One recovers the action of fluctuating hydrodynamics

[Spohn, Bertini De Sole Gabrielli Jona-Lasinio Landim]

$\psi(s)$: again[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

Periodic boundary conditions

More general fluctuating hydrodynamics

$$\frac{1}{Lt} \langle Q \rangle \propto D(\rho) \quad (\text{Fourier's law})$$

$$\frac{1}{Lt} \langle Q^2 \rangle_c = \sigma(\rho) \quad (\text{For the SSEP, } \sigma = \rho(1 - \rho))$$

$\psi(s)$: again

[Appert, Derrida, VL, van Wijland, PRE 78 021122]

Periodic boundary conditions

More general fluctuating hydrodynamics

$$\frac{1}{Lt} \langle Q \rangle \propto D(\rho) \quad (\text{Fourier's law})$$

$$\frac{1}{Lt} \langle Q^2 \rangle_c = \sigma(\rho) \quad (\text{For the SSEP, } \sigma = \rho(1 - \rho))$$

Saddle point evaluation

$$\langle e^{-sQ} \rangle \sim \int \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L S_s[\hat{\rho}, \rho]\}$$

$\psi(\mathbf{s})$: again

[Appert, Derrida, VL, van Wijland, PRE 78 021122]

Periodic boundary conditions

More general fluctuating hydrodynamics

$$\frac{1}{Lt} \langle Q \rangle \propto D(\rho) \quad (\text{Fourier's law})$$

$$\frac{1}{Lt} \langle Q^2 \rangle_c = \sigma(\rho) \quad (\text{For the SSEP, } \sigma = \rho(1 - \rho))$$

Saddle point evaluation

$$\langle e^{-sQ} \rangle \sim \int D\rho D\hat{\rho} \exp\{L S_s[\hat{\rho}, \rho]\}$$

Large deviation function

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{2L} \mathbf{s}^2 \frac{\langle Q^2 \rangle_c}{t}}_{\text{at saddle-point}} + \underbrace{L^{-2} D\mathcal{F}(u)}_{\substack{\text{small fluctuations} \\ \text{(determinant)}}} \quad \text{with} \quad u = -\mathbf{s}^2 \frac{\sigma\sigma''}{16D^2}$$

Correspondence between
(Gaussian) integration of small fluctuations
AND
discreteness of Bethe root repartition.

More general?

Correspondence between
(Gaussian) integration of small fluctuations
AND
discreteness of Bethe root repartition.

More general?

Fluctuating hydrodynamics for quantum chains?

With a field

[Appert, Derrida, VL, van Wijland, PRE **78** 021122]

Periodic boundary conditions

Driving field E

With a field

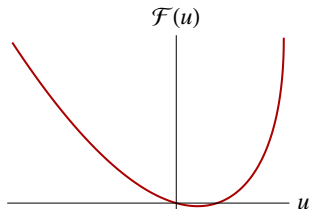
[Appert, Derrida, VL, van Wijland, PRE 78 021122]

Periodic boundary conditions

Driving field E

Large deviation function

$$\psi(s) = \underbrace{\frac{1}{2L} s(s-E) \frac{\langle Q^2 \rangle_c}{t}}_{\text{at saddle-point}} + \underbrace{L^{-2} D\mathcal{F}(u)}_{\text{small fluctuations (determinant)}} \quad \text{with} \quad u = \underbrace{-s(s-E) \frac{\sigma\sigma''}{16D^2}}_{\text{can become } > 0}$$



Dynamical phase transition
between
stationary and non-stationary
profiles

Outline

- 1 Motivations
- 2 Microscopic approach
 - Operator approach
 - Bethe Ansatz
- 3 Macroscopic approach
 - Fluctuating hydrodynamics
 - Finite-size corrections
- 4 Mapping non-equilibrium to equilibrium
 - For the integrated current
 - For the density profile



Microscopic approach [Imparato, VL, van Wijland, arXiv:0911.0564]

Large deviations of the current

$$\psi(\mathbf{s}) = \max_{\text{Sp}} \mathbb{W}(\mathbf{s})$$

$$\mathbb{W}(\mathbf{s}) = \overbrace{\sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1}}^{\text{invariant by rotation}}$$

$$+ \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1]$$

$$+ \delta [S_L^+ e^{\mathbf{s}} - (1 - \hat{n}_L)] + \beta [S_L^- e^{-\mathbf{s}} - \hat{n}_L]$$

Microscopic approach [Imparato, VL, van Wijland, arXiv:0911.0564]

Large deviations of the current

$$\psi(\mathbf{s}) = \max_{\text{Sp}} \mathbb{W}(\mathbf{s})$$

$$\mathbb{W}(\mathbf{s}) = \overbrace{\sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1}}^{\text{invariant by rotation}} + \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1] + \delta [S_L^+ e^{\mathbf{s}} - (1 - \hat{n}_L)] + \beta [S_L^- e^{-\mathbf{s}} - \hat{n}_L]$$

Local transformation

$$\mathcal{Q}^{-1} \mathbb{W}(\mathbf{s}) \mathcal{Q} = \sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1} + \alpha' [S_1^+ - (1 - \hat{n}_1)] + \gamma' [S_1^- - \hat{n}_1] + \delta' [S_L^+ e^{\mathbf{s}'} - (1 - \hat{n}_L)] + \beta' [S_L^- e^{-\mathbf{s}'} - \hat{n}_L]$$

describes contact with reservoirs of same densities

Microscopic approach [Imparato, VL, van Wijland, arXiv:0911.0564]

Large deviations of the current

$$\psi(\mathbf{s}) = \max_{\text{Sp}} \mathbb{W}(\mathbf{s})$$

$$\mathbb{W}(\mathbf{s}) = \overbrace{\sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1}}^{\text{invariant by rotation}} + \alpha [S_1^+ - (1 - \hat{n}_1)] + \gamma [S_1^- - \hat{n}_1] + \delta [S_L^+ e^{\mathbf{s}} - (1 - \hat{n}_L)] + \beta [S_L^- e^{-\mathbf{s}} - \hat{n}_L]$$

Local transformation **Fluctuations@eq** \leftrightarrow **Fluctuations@non-eq**

$$\mathcal{Q}^{-1} \mathbb{W}(\mathbf{s}) \mathcal{Q} = \sum_{1 \leq k \leq L-1} \vec{S}_k \cdot \vec{S}_{k+1} + \alpha' [S_1^+ - (1 - \hat{n}_1)] + \gamma' [S_1^- - \hat{n}_1] + \delta' [S_L^+ e^{\mathbf{s}'} - (1 - \hat{n}_L)] + \beta' [S_L^- e^{-\mathbf{s}'} - \hat{n}_L]$$

describes contact with reservoirs of same densities

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

$$\langle e^{-sQ} \rangle \sim \int \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L\mathcal{S}_s[\hat{\rho}, \rho]\}$$

Fluctuations $\phi, \hat{\phi}$ around the saddle

$$\rho(\mathbf{x}, t) = \rho_c(\mathbf{x}) + \phi(\mathbf{x}, t)$$

$$\hat{\rho}(\mathbf{x}, t) = \hat{\rho}_c(\mathbf{x}) + \hat{\phi}(\mathbf{x}, t)$$

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

$$\langle e^{-sQ} \rangle \sim \int \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L\mathcal{S}_s[\hat{\rho}, \rho]\}$$

Fluctuations $\phi, \hat{\phi}$ around the saddle

$$\rho(\mathbf{x}, t) = \rho_c(\mathbf{x}) + \phi(\mathbf{x}, t) \qquad \hat{\rho}(\mathbf{x}, t) = \hat{\rho}_c(\mathbf{x}) + \hat{\phi}(\mathbf{x}, t)$$

Mapping of non-eq. fluctuations $\phi, \hat{\phi}$ to eq. fluctuations $\phi', \hat{\phi}'$

$$\begin{aligned} \phi(\mathbf{x}, t) &= (\partial_x \hat{\rho}_c)^{-1} \phi'(\mathbf{x}, t) + (\partial_x \rho_c)^{-1} \hat{\phi}'(\mathbf{x}, t) \\ \hat{\phi}(\mathbf{x}, t) &= (\partial_x \hat{\rho}_c) \hat{\phi}'(\mathbf{x}, t) \end{aligned}$$

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

$$\langle e^{-sQ} \rangle \sim \int \mathcal{D}\rho \mathcal{D}\hat{\rho} \exp\{L \mathcal{S}_s[\hat{\rho}, \rho]\}$$

Fluctuations $\phi, \hat{\phi}$ around the saddle

$$\rho(\mathbf{x}, t) = \rho_c(\mathbf{x}) + \phi(\mathbf{x}, t) \qquad \hat{\rho}(\mathbf{x}, t) = \hat{\rho}_c(\mathbf{x}) + \hat{\phi}(\mathbf{x}, t)$$

Mapping of non-eq. fluctuations $\phi, \hat{\phi}$ to eq. fluctuations $\phi', \hat{\phi}'$

$$\begin{aligned} \phi(\mathbf{x}, t) &= (\partial_x \hat{\rho}_c)^{-1} \phi'(\mathbf{x}, t) + (\partial_x \rho_c)^{-1} \hat{\phi}'(\mathbf{x}, t) \\ \hat{\phi}(\mathbf{x}, t) &= (\partial_x \hat{\rho}_c) \hat{\phi}'(\mathbf{x}, t) \end{aligned}$$

Large deviation function

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L} \mu(\mathbf{s})}_{\text{saddle}} + \underbrace{\frac{D}{8L^2} \mathcal{F} \left(\frac{\sigma''}{2D^2} \mu(\mathbf{s}) \right)}_{\text{fluctuations}} \quad \text{same } \mathcal{F} \text{ as at eq.}$$

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

Physical remark

In contact with reservoirs, neither SSEP nor KMP present a 'dynamical phase transition' (\Leftarrow no singularity in $\psi(s)$).

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

Physical remark

In contact with reservoirs, neither SSEP nor KMP present a 'dynamical phase transition' (\Leftarrow no singularity in $\psi(\mathbf{s})$).

Technical remark

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L}\mu(\mathbf{s})}_{\text{saddle}} + \underbrace{\frac{D}{8L^2}\mathcal{F}\left(\frac{\sigma''}{2D^2}\mu(\mathbf{s})\right)}_{\text{fluctuations}}$$

the **saddle** also contributes to the order $1/L^2$

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

Physical remark

In contact with reservoirs, neither SSEP nor KMP present a 'dynamical phase transition' (\Leftarrow no singularity in $\psi(\mathbf{s})$).

Technical remark

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L}\mu(\mathbf{s})}_{\text{saddle}} + \underbrace{\frac{D}{8L^2}\mathcal{F}\left(\frac{\sigma''}{2D^2}\mu(\mathbf{s})\right)}_{\text{fluctuations}}$$

the **saddle** also contributes to the order $1/L^2$

one way to compute these finite-size corrections to **saddle**:

- start from $\mathbb{W}(\mathbf{s})$
- use coherent states to obtain a *spatially-discrete* action $S_s[\rho, \hat{\rho}]$
- solve the *spatially-discrete* stationary saddle-point equations

Macroscopic approach

[Imparato, VL, van Wijland, **PRE** 80 011131]

Physical remark

In contact with reservoirs, neither SSEP nor KMP present a 'dynamical phase transition' (\Leftarrow no singularity in $\psi(\mathbf{s})$).

Technical remark

$$\psi(\mathbf{s}) = \underbrace{\frac{1}{L}\mu(\mathbf{s})}_{\text{saddle}} + \underbrace{\frac{D}{8L^2}\mathcal{F}\left(\frac{\sigma''}{2D^2}\mu(\mathbf{s})\right)}_{\text{fluctuations}}$$

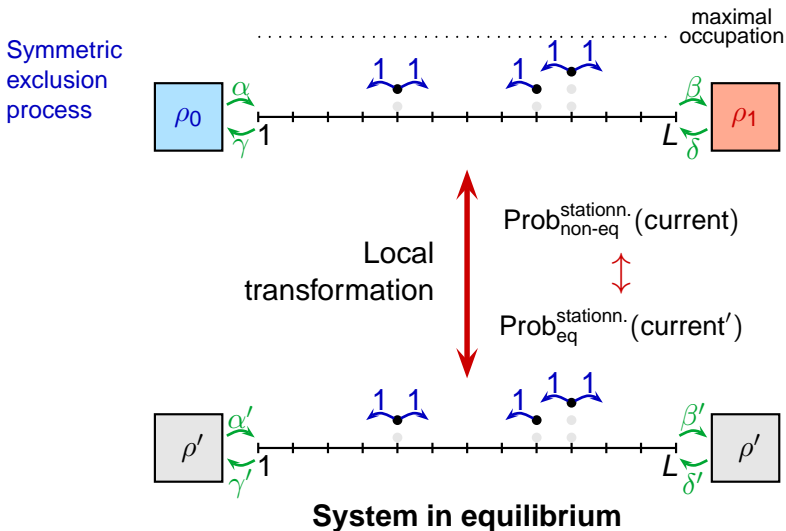
the **saddle** also contributes to the order $1/L^2$

one way to compute these finite-size corrections to **saddle**:

- start from $\mathbb{W}(\mathbf{s})$
- use coherent states to obtain a *spatially-discrete* action $S_s[\rho, \hat{\rho}]$
- solve the *spatially-discrete* stationary saddle-point equations

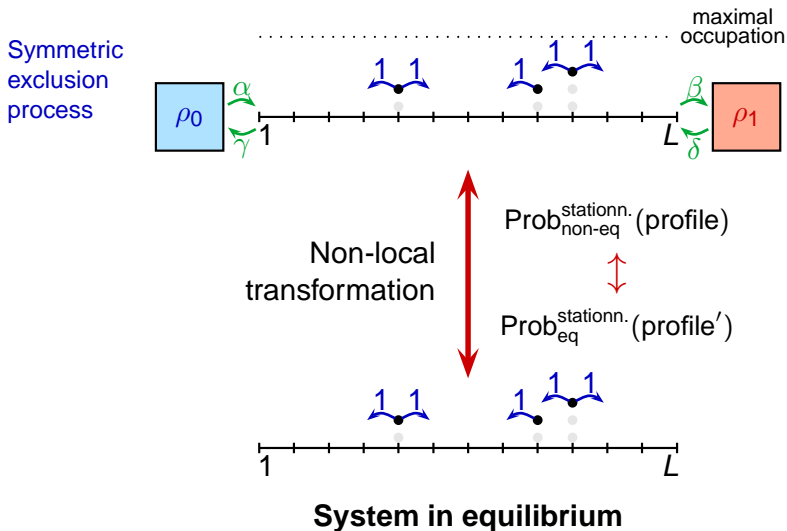
→ How does this translate in the Bethe Ansatz approach?

For the current

[Imparato, VL, van Wijland, **PRE** 80 011131]

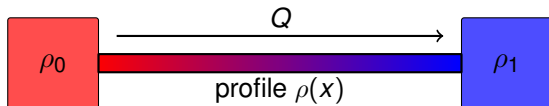
For the density profile

[Tailleur, Kurchan, VL, JPA 41 505001]



Non-local mapping

[Tailleur, Kurchan, VL, JPA **41** 505001]

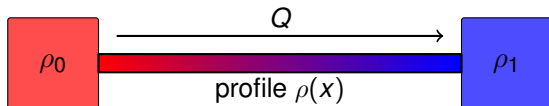


Boundary-driven transport model:

- long-range correlations
- breaking of time-reversal symmetry

Non-local mapping

[Tailleur, Kurchan, VL, JPA **41** 505001]



Boundary-driven transport model:

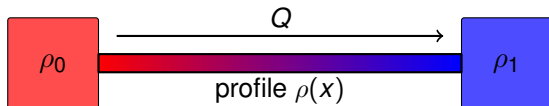
- long-range correlations
- breaking of time-reversal symmetry

Non-local mapping to equilibrium:

- accounts for long-range correlations
- (density gradient)_{non-eq.} \longleftrightarrow (fixed density)_{eq.}
- yields $\text{Prob}[\rho(x)]$ through a maximisation principle

Non-local mapping

[Tailleur, Kurchan, VL, JPA **41** 505001]



Boundary-driven transport model:

- long-range correlations
- breaking of time-reversal symmetry

Non-local mapping to equilibrium:

- accounts for long-range correlations
- (density gradient)_{non-eq.} \longleftrightarrow (fixed density)_{eq.}
- yields $\text{Prob}[\rho(x)]$ through a maximisation principle

→ Applies to non-equilibrium quantum chains?

Summary

Approach:

- operator formalism
- large deviation function

Extensions:

- fluctuating hydrodynamics
- saddle-point method, instantons
- integration of fluctuations (dynamical phase transition)

Open questions:

- Eq \leftrightarrow non-eq mapping in higher dimensions?
- More generic systems of interacting particles?
- Link btw the eq \leftrightarrow non-eq mappings?
- Crossover to KPZ? Other universal fluctuations?